

Simulation of Hydrophobicity Evaluation and Structural Optimization Design Method for Micro-Array Units

Yang Shen^{1*}, Ge Guo², Yafeng Zhang^{3*}, Yuan Lu¹, Xin Wang¹

¹Army Artillery Air Defense Academy, Chinese People's Liberation Army, Hefei, China

²Department of Border Defence Command, Army Academy of Border and Coastal Defence in Urumqi, Urumqi, China

³Department of Equipment Support, Automobile NCO College, Army Military Transportation University, Bengbu, China
Email: *willnudit@sina.com, solosad.ok@163.com

How to cite this paper: Shen, Y., Guo, G., Zhang, Y.F., Lu, Y. and Wang, X. (2018) Simulation of Hydrophobicity Evaluation and Structural Optimization Design Method for Micro-Array Units. *Journal of Surface Engineered Materials and Advanced Technology*, 8, 37-48.

<https://doi.org/10.4236/jsemat.2018.82004>

Received: March 28, 2018

Accepted: April 22, 2018

Published: April 25, 2018

Copyright © 2018 by authors and Scientific Research Publishing Inc. This work is licensed under the Creative Commons Attribution International License (CC BY 4.0).

<http://creativecommons.org/licenses/by/4.0/>



Open Access

Abstract

In recent years, researches published on hydrophobic materials increase rapidly, wherein the method for changing hydrophobicity by modifying a micro-array structure on the surface of the material also has been proposed. Of course, if it is possible for us to quantitatively analyse and evaluate hydrophobicity of different structures of one certain material at first, this task will greatly optimize the design of actual structures. In this work, we used the algorithm for Laminar Two-Phase Flow, Horizontal-set method integrated in COMSOL to build two single-pore simulation structures in different shapes and simulated the behaviour of the liquid transition from Cassie-state to Wenzel-state during the impregnation process. After that, the intrinsic contact angle of Structure T (a porous structure with a T-shaped sectional profile) was obtained under a certain pressure which maintained liquids in Cassie-state. Meanwhile, two equilibrium states of the liquid-air interface as well as two different patterns of the equilibrium state disrupting were found in Structure R (a porous structure with a Chamfered T-shaped sectional profile). Simulation results show that the modelling method can be applied for simulating the hydrophobicity of different porous structures and optimizing the procedures for the design of the micro-array efficiently.

Keywords

Hydrophobic Materials, Micro-Array, Laminar Two-Phase Flow, Horizontal-Set Method, COMSOL

1. Introduction

Hydrophobic materials have become widely used due to their advantages of self-cleaning, resistance reduction and anti-icing etc. [1]. To be specific, the use of hydrophobic materials in the appropriate part of the architectural structure can reduce the corrosion from the acid rain: as in ships, it can reduce the resistance during the navigation; also as in pipelines, it can effectively reduce the corrosion; thereby improving the transportation efficiently. To obtain a material with the better hydrophobic property, one method which has been studied and adopted is to manufacture a micro-array on the surface.

The design idea of the hydrophobic surface micro-array is adapted from the natural existence of the hydrophobic surface while a variety of hydrophobic material surfaces have been designed and manufactured. To study the behaviour of droplets spreading on a solid surface, Cassie and Wenzel firstly classified the infiltration state into Cassie-state and Wenzel-state. Cassie-state is an ideal hydrophobic state that the liquid is sheared up by a rough surface micro-structure [2]; nevertheless, the liquid can not fill up the pit on the surface but form compound contact because of the air trapped with the surface (shown in **Figure 1(a)**). Wenzel-state is that the liquid can completely fill up a rough surface while the actual contact area is larger than the apparent contact area (shown in **Figure 1(b)**) [3].

Scholars have scientifically explained the hydrophobicity of various micro-structures through mathematical modelling, which made hydrophobic theory continuously improve. B. Emami *et al.* proved that the cylindrical pore maximized the critical pressure [4]. Rene Hensel *et al.* fitted the microscopic structure on the surface of the Springtails skin and calculated the surface infiltration of the liquid [5]. In addition, they proved that both the Mushroom-shaped cross-sectional pore and the Chamfered T-shaped cross-sectional hole had excellent hydrophobicity. By transforming the concave structure of a Chamfered T-shaped hole into a convex structure and basing on a micro-fabrication process, Tingyi Liu *et al.* obtained an array of surface micro-structures, which made the hydrophilic silica surface exhibit excellent super-hydrophobic characteristics [6]. Takahiro Koishi *et al.* simulated the infiltration behaviour of water droplets composed of a large number of water molecules on the surface of a

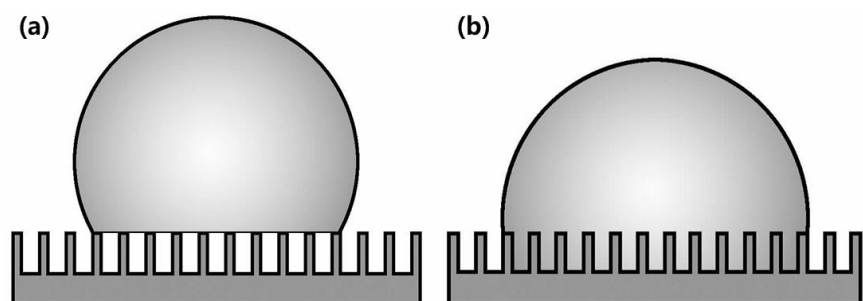


Figure 1. (a) Cassie-state; (b) Wenzel-state.

modified micro-array based on molecular dynamics simulations [7]. Jie Xiao *et al.* simulated the adhesion behaviour of water droplets on the surface of the modified micro-array when there was wind blowing [8]. It could be said that further improvement of theories mentioned above could provide an important reference for the design of hydrophobic micro-array.

Apparent contact angle is a macroscopic representation of the droplet infiltration behaviour influenced by intrinsic contact angle near each unit of the array, which can reflect hydrophobic of the material. Therefore, it can provide guiding significance for grasping the whole retention of droplets on the surface and a basis for the surface design. Necessarily, to tracing the liquid-gas transitional surface, we used Horizontal-set method integrated in COMSOL to define horizontal-set function φ (always regarded as the volume fraction of the liquid phase in the transitional region) transferring from the liquid phase ($\varphi = 1$) to the gas phase ($\varphi = 0$). The gas-liquid transitional surface can be calculated by Equation (1) [9].

$$\frac{\partial \varphi}{\partial t} + \mathbf{u} \cdot \nabla \varphi = \gamma \mathcal{N} \cdot \left[\varepsilon \nabla \varphi - \varphi(1-\varphi) \frac{\nabla \varphi}{|\nabla \varphi|} \right] \quad (1)$$

In the equation, ε represents the thickness of defined transition surface; γ represents reinitialization parameter whose value is generally set as the maximum flow velocity of the liquid; \mathbf{u} represents the fluid velocity.

Figure 2 shows a surface micro-pore model of the array applied for the simulation analysis. It can be seen that the plane of the model is a concave structure; actually, the structure is prepared on the whole surface of the materials generally with the nether base closed. When sufficient amount of fluid a minority of micro-units are covered under sufficient amount of fluid, the pressure in the pore will go larger as the increase of the fluid entering, which generally makes the material hydrophobic. However, in order to provide some convenience for the design of the non-pore convex array, the nether base of the pore will not be closed but connect to the air (shown in **Figure 2(b)**) and the structural profile is T-shaped (shown in **Figure 2(c)**). When the droplet is on the surface, the final stability of droplets will be affected by its gravity, the gas pressure in the pore, and the infiltration of the solid-liquid surface. If droplets are finally in Cassie-state (shown in **Figure 3(a)**), the liquid will flow as **Figure 3(b)** and **Figure 3(c)** show. When the liquid enters into the cylindrical pore and the liquid gravity, capillary force, Young's pressure, and air pressure in the pore all reaches an equilibrium state, the liquid will form stable droplets in the pore. The other section profile in this paper is Structure R, of which the nether surface is not closed.

2. Simulation Details

The design dimension of Structure T is shown in **Figure 4**. The whole surface exits an extension of the pore of which the length L is 10 μm , the thickness t is 1 μm and the distance B (from the lateral wall of the nether pore to the central axis) is 55 μm . No sooner were L and B determined than the reference axis

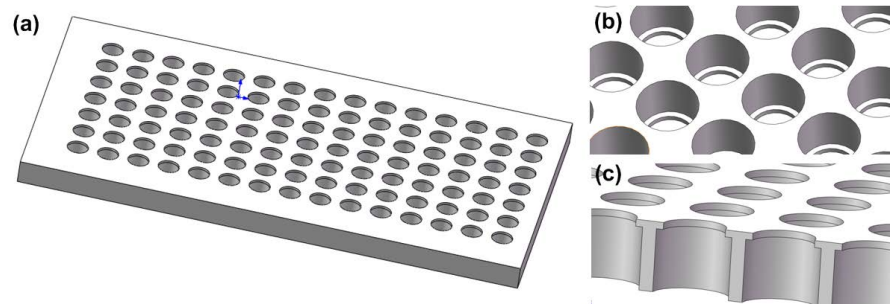


Figure 2. (a) Top view of the surface of the array; (b) Partial enlarged drawing of the back of the array; (c) Sectional drawing of the structure.

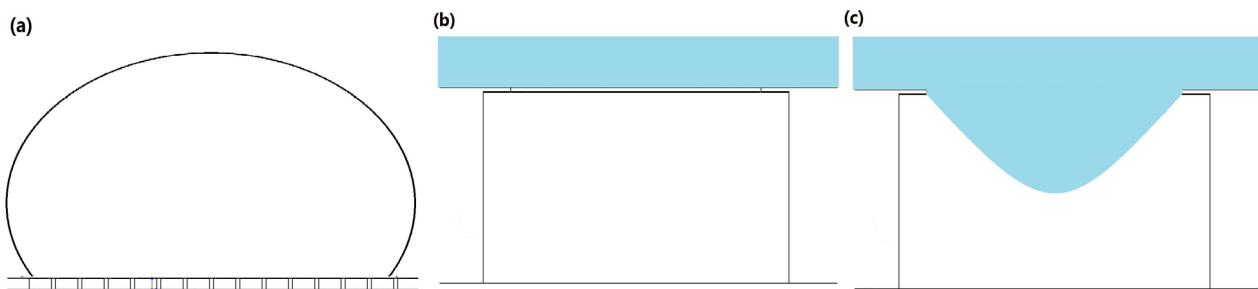


Figure 3. (a) Shape of the droplet in Cassie-state; (b) Original state between the droplet and the pore; (c) Stable state after the droplet entering into the pore.

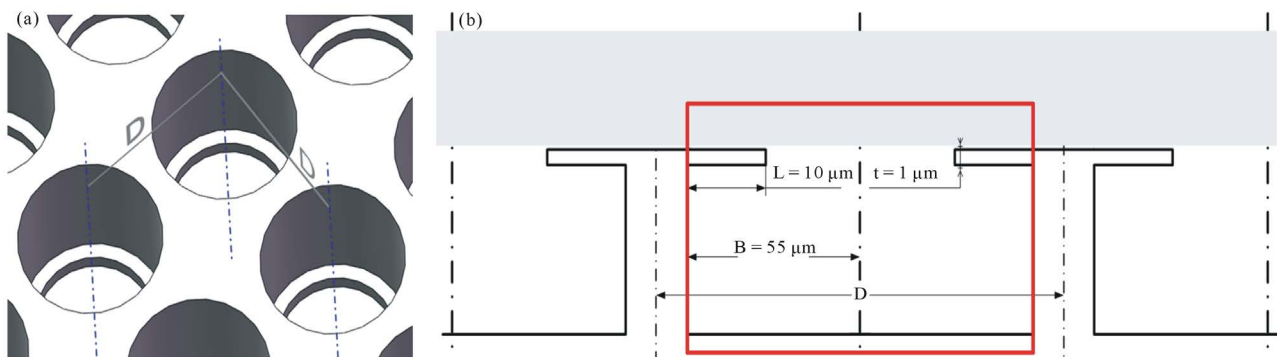


Figure 4. (a) View of the nether surface of the array; (b) Sizes of the microarray unit.

distance D , of which the macroscopic appearance is the variation of apparent contact angle of droplets, affected not only the contact area with droplets but also the pressure acted on them. As we only performed quantitative analysis by the time we set the inlet pressure, we did not need to calculate the value of apparent contact angle. Meanwhile, we just built parts of the model in the red box (shown in **Figure 4(b)**). Thus, D would be out of considered. According to the whole progress, the model could establish a rotational symmetric one to simplify the calculation. Additionally, as the solid material was only connected with the air and the liquid, we merely needed to define the infiltrating wall and stimulate from the origin state to the end of the pore fulfilling with phase initialization despite modeling the entire solid phase.

2.1. The establishment of the Structure T model

Figure 5 and **Figure 6** separately shows the rotating symmetrical three-dimensional model and the two-dimensional model (inflexion-point 1 marked at the end of the structure) of one unit of the array. We defined the target liquid as water and different contact angles of the humid wall to represent for different materials at first. Then, if only considering the gravity of droplets and the inlet pressure order of $10^2/\text{Pa}$, we defined the water pressure at the inlet; However, we observed that the order was so small as to make the result in conformity with the present literature, which further explained merely considering the gravity was not accurate. So we set the value to a larger order of $10^2/\text{Pa}$ at last [5]. When the apparent contact angle was larger than 90° , the droplet was a fixed volume.

When applying Horizontal-set method to analyse the interface between the liquid and the gas, we defined the thickness of the liquid-gas transition surface ε as 5×10^{-7} m. As the thickness value of the liquid-gas transition surface was larger, whether inflexion-point 1 is the right angle or the round angle had no difference [10]. Considering the manufacturing process of the structure while to observe three-phase contact line, we set inflexion-point 1 as a chamfer of which the radius was $0.2 \mu\text{m}$.

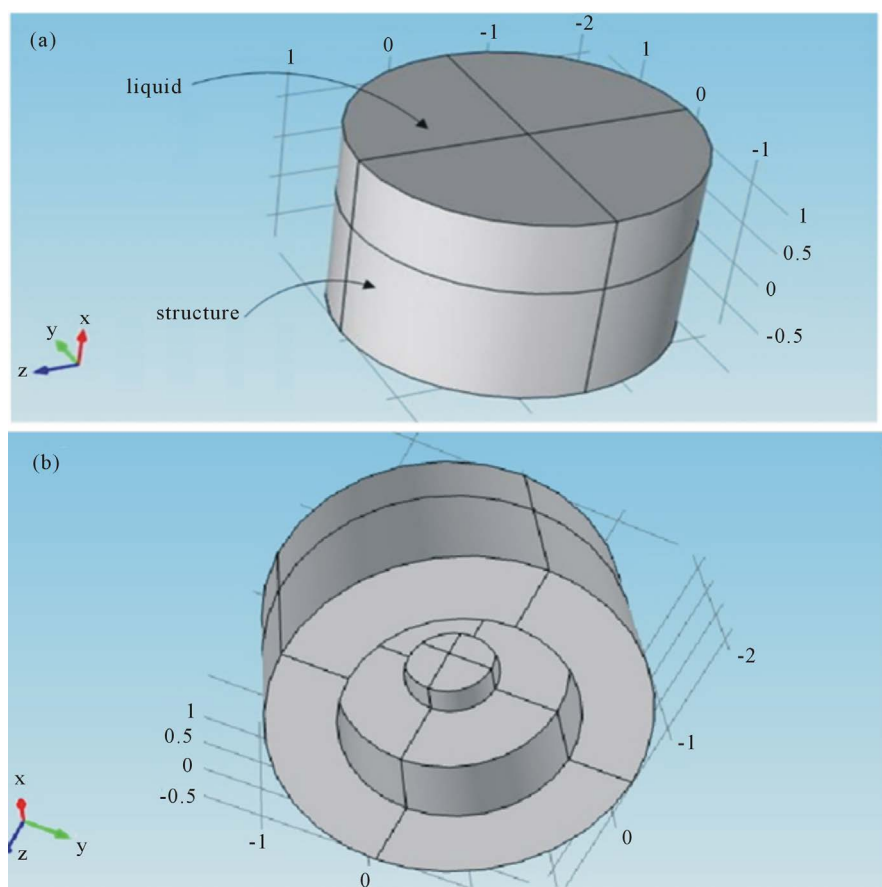


Figure 5. (a) Top view of the model of structure T; (b) Top view of the model of structure f.

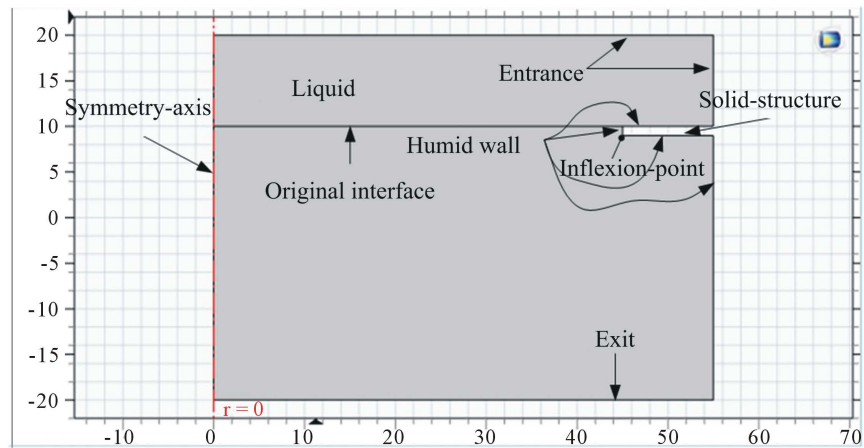


Figure 6. 2-D view of the model.

2.2. The Establishment of Structure R Model

Figure 7(a) and **Figure 7(b)** separately shows the rotating symmetrical three-dimensional model and the two-dimensional model. The model establishment process is similar to Structure T. The angular thickness of Structure R t equals $1\ \mu\text{m}$, the angular height h equals $4\ \mu\text{m}$ and the length of the structure edge L equals $10\ \mu\text{m}$. The nether left corner is marked as inflexion-point 1, and the nether right corner is marked as inflexion-point 2.

3. Results and Discussions

3.1. The Results of Structure T Simulation

In the Structure T simulation model, we set the value of the inlet boundary pressure to be $1000\ \text{Pa}$ and gradually increased intrinsic contact angle by the step distance of 0.5° . The simulation results showed that three-phase contact line would firstly cross inflexion-point 1, after some reciprocating motion, it would finally keep balance at inflexion-point 1 when the intrinsic contact angle equalled 18.36° (shown in **Figure 8**, where blue represents the liquid and red represents the air).

When intrinsic contact angle was decreased to 17.86° , the origin inlet boundary pressure would be larger so that the droplet could not form Cassie-state on the surface at all. Thus, three-phase contact line would firstly go forward to a longer distance and then backwards to inflexion-point 1; after that, it would move to the right boundary of the nether surface and continue to move. This implies three-phase contact line moved to the right boundary, thus the liquid-gas interface was unable to achieve an equilibrium state (shown in **Figure 9**, where blue represented the liquid and red represented the air).

When intrinsic contact angle increased to 18.86° , the origin inlet boundary pressure would smaller so that the structure of the material would be more difficult to infiltrate. The infiltrating behaviour of three-phase contact line would be similar to the infiltrating behaviour in condition that the intrinsic original

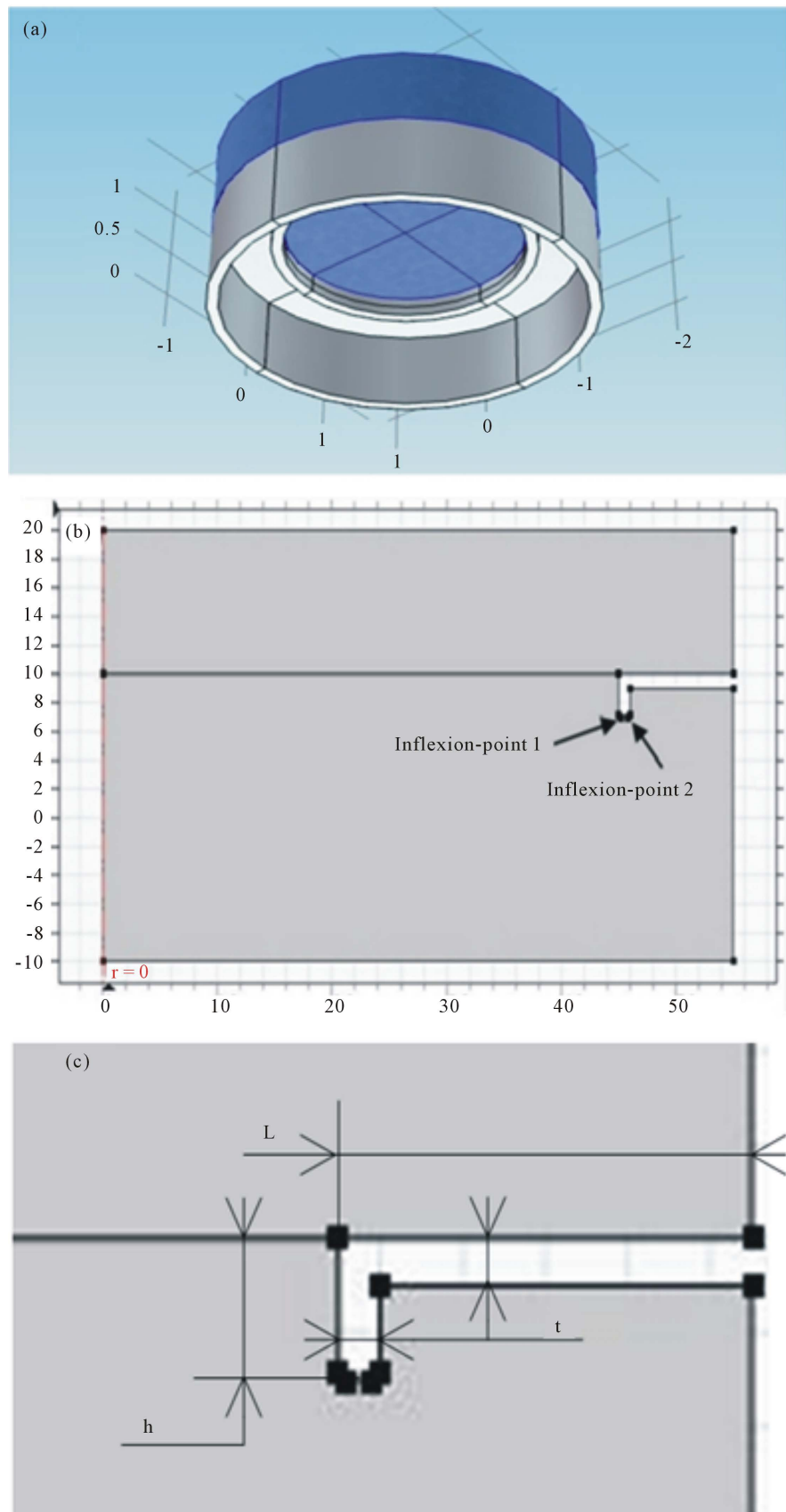


Figure 7. (a) Model of structure R; (b) 2-D view of the model; (c) Partial enlarged view.

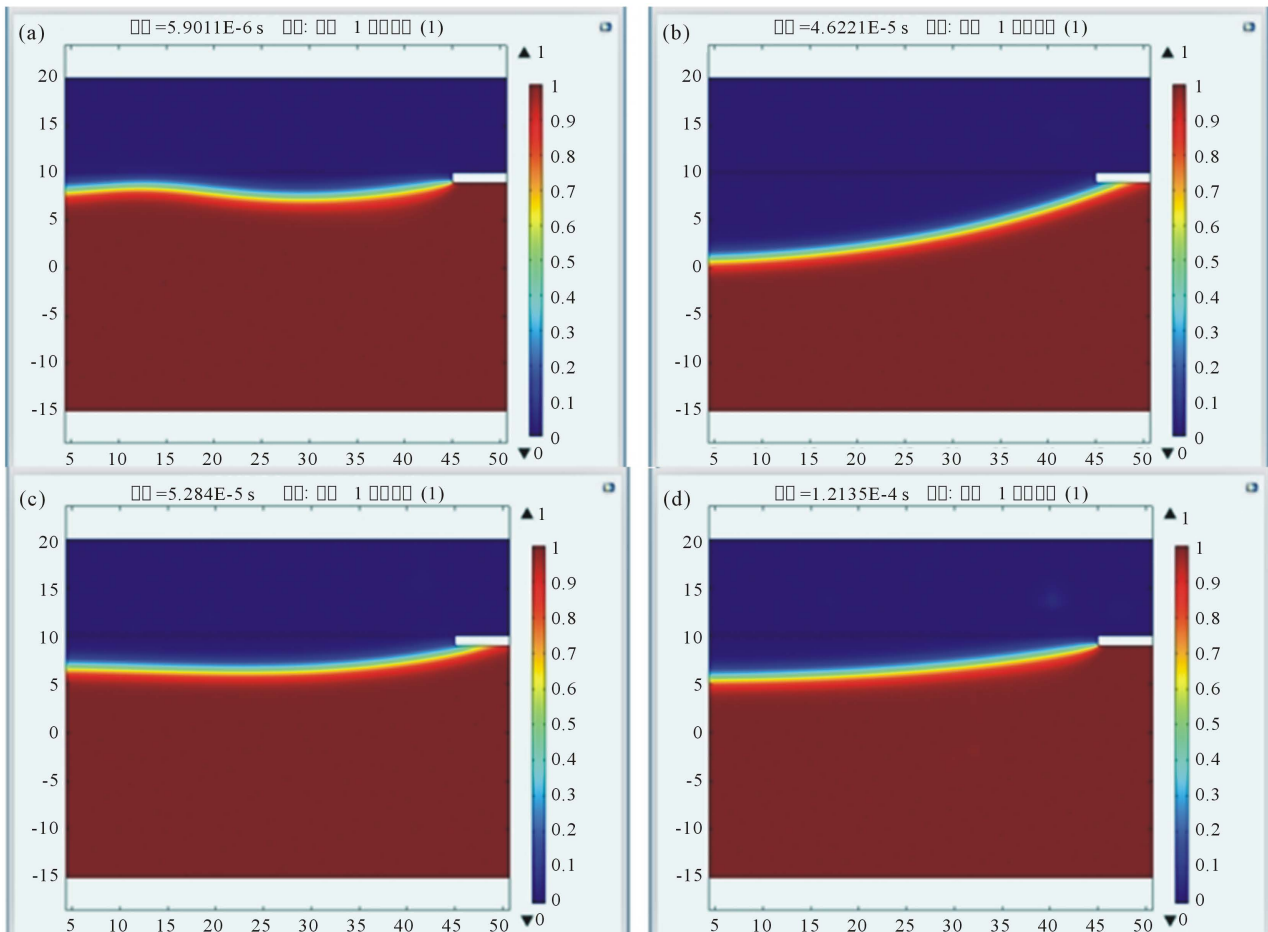


Figure 8. Oscillation of three-phase contact line near inflexion-point 1 at the angle of 18.36° : (a) Time of 5.9×10^{-6} s; (b) Time of 54.62×10^{-5} s; (c) Time of 5.28×10^{-5} s; (d) Time of 1.21×10^{-4} s.

contact angle equalled 18.36° ; Also, its humid distance would be shorter after going cross inflexion-point 1 (shown in **Figure 10**, where blue represented the liquid and red represented the air).

3.2. The Results of Structure R Simulation

It could be seen that the equilibrium state at inflexion-point 1 was similar to that of Structure T in spite of different increasing thicknesses. As for inflexion-point 2, the inlet pressure was set to 1000 Pa and intrinsic contact angle of the humid wall was 17.86° . Owing to the thickness of the angle was smaller than the humid distance of three-phase contact line passing over inflexion-point 1 in Structure T, we observed three-phase contact line moved to inflexion-point 2 (not passed over it) then returned to inflexion-point 1; After that, it would not sustain oscillation until balance. Continuously reducing contact angle to 1° and 0.01° , we observed the equilibrium state near inflexion-point 1 (shown in **Figure 11** and **Figure 12**).

Increasing the inlet pressure to 2000 Pa, 2500 Pa and 3000 Pa, which meant to change apparent contact angle of the droplet, we obtained the simulation results

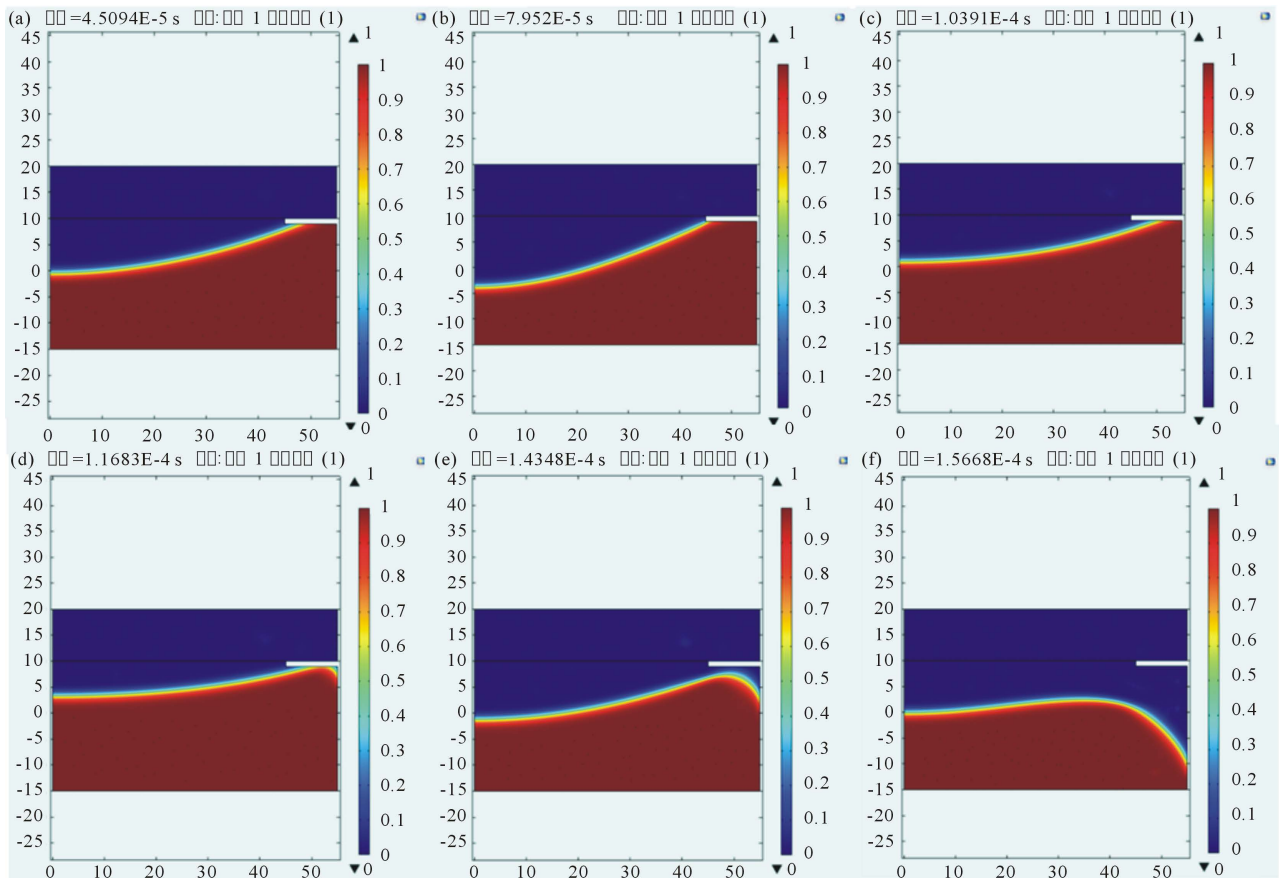


Figure 9. Motion of three-phase contact line from inflexion-point 1 at the angle of 17.86°: (a) Time of 4.5×10^{-5} s; (b) Time of 7.95×10^{-5} s; (c) Time of 1.03×10^{-4} s; (d) Time of 1.17×10^{-4} s; (e) Time of 1.43×10^{-4} s; (f) Time of 1.57×10^{-4} s.

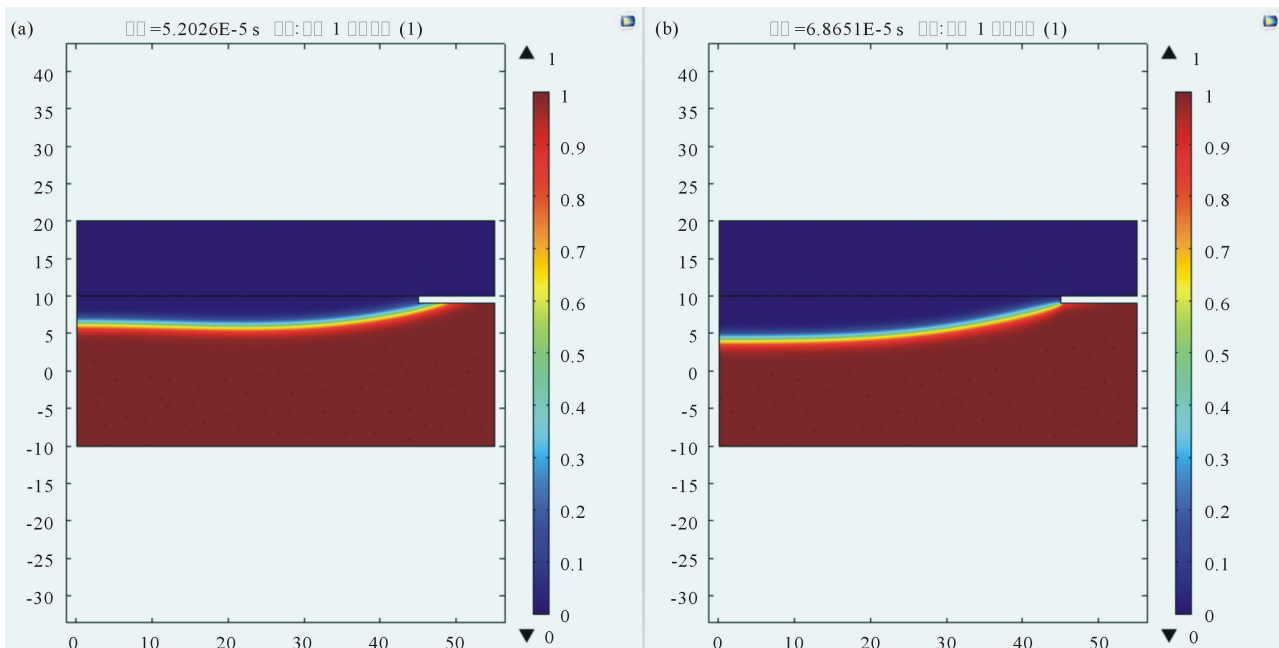


Figure 10. Motion of three-phase contact line from inflexion-point 1 at the angle of 18.86°: (a) Time of 5.203×10^{-5} s; (b) Time of 6.865×10^{-5} s.

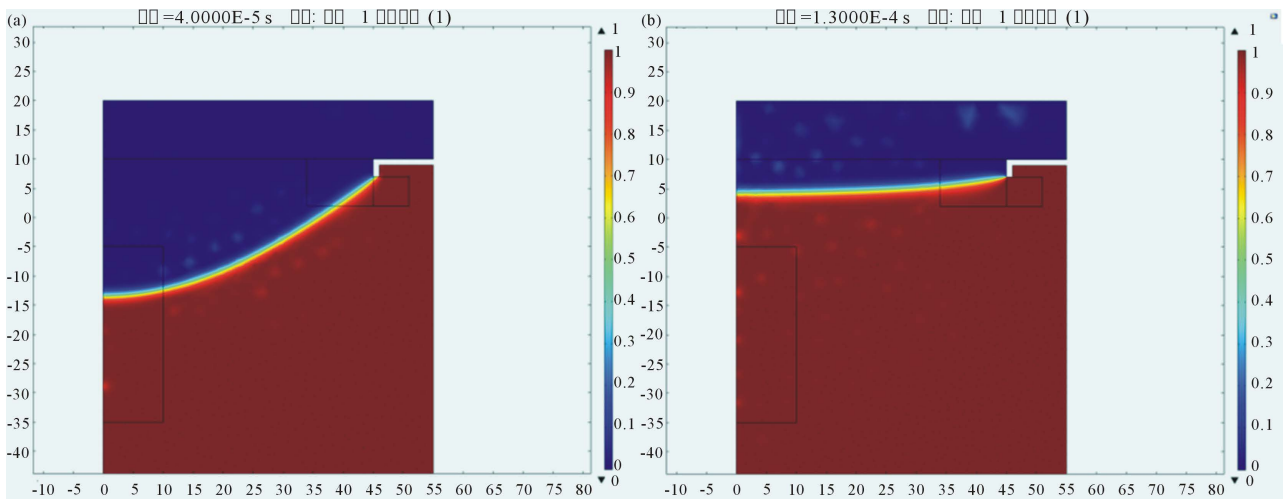


Figure 11. Motion of three-phase contact line from inflexion-point 1 at the angle of 1° : (a) Time of 4.0×10^{-5} s; (b) Time of 1.03×10^{-4} s.

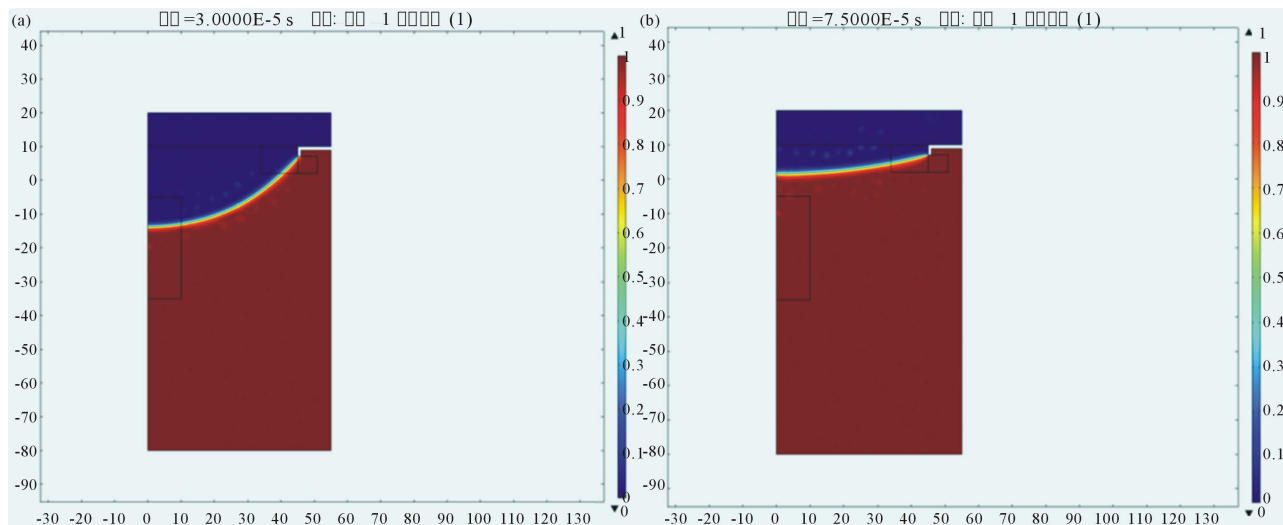


Figure 12. Motion of three-phase contact line from inflexion-point 1 at the angle of 0.01° : (a) Time of 3.0×10^{-5} s; (b) Time of 7.5×10^{-5} s.

(shown in **Figures 13(a)-(c)**). When the inlet pressure reached to 2000 Pa, three-phase contact line could not cross inflexion-point 2 because of the gravity of the droplet, but returned to inflexion-point 1 then achieved an equilibrium state. When the inlet pressure reached to 2500 Pa, three-phase contact line would go across inflexion-point 2, and moved upward along with the inner wall of the chamfer, but could not achieve the equilibrium state. When the inlet pressure reached to 3000 Pa, the interface of the liquid-gas would contact the structural side wall and could not achieve the equilibrium state because of the large radius of the droplets.

4. Conclusions

- 1) From Structure T simulation conclusion, we proved the existence of the

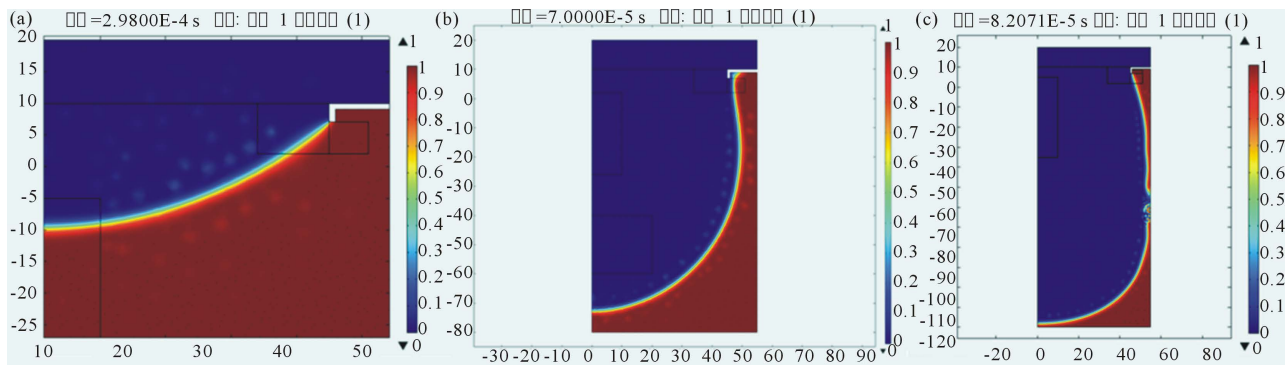


Figure 13. Different states under pressures of 2000 Pa, 2500 Pa and 3000 Pa.

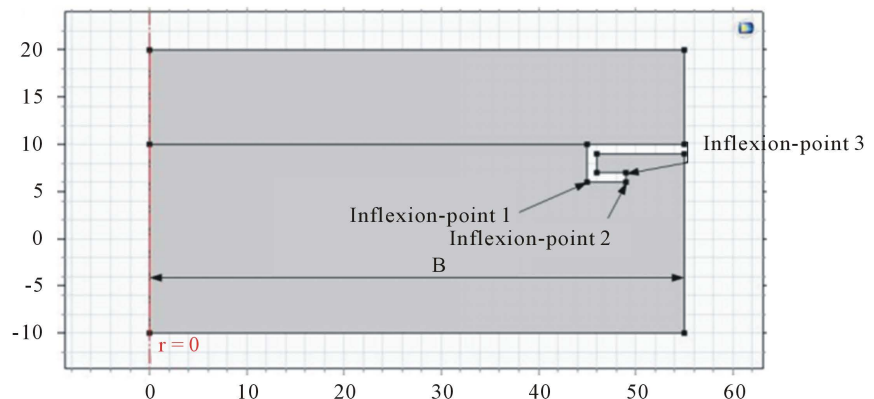


Figure 14. 2-D model for hydrophobicity comparison.

critical intrinsic contact angle by the inlet pressure valuing 18.36° . Another, the inlet pressure needs to reduce with the decreasing of intrinsic contact angle.

2) As for the same material, the inlet pressure of Structure R whose liquid-gas interface reached to an equilibrium state was much larger than that of Structure T [6].

5. Future Prospects

Because the macroscopic wettability is affected by the solid/air ratio of the droplet in contact with the engineered solid surface, we will further consider the spacing between adjacent pores and make clear how the intrinsic contact and surface tension are formulated.

In the design of the pore structure size, we can further increase the radius of the nether pore and reduce the radius of the upper pore to make droplets inside not contact the pore wall. Also, we can verify whether some similar structures were hydrophobic or not by changing the shape. For example, as **Figure 14** shows, we can judge the hydrophobicity of the material from whether three-phase contact line can go across inflexion-point 3 or not when the inlet pressure is 2500 Pa. In other hand, it can be necessary to increase the distance between the lateral wall of the nether pore and the center of B when compared with **Figure 7(b)**.

References

- [1] Feng, L., Li, S. and Li, Y. (2002) Super-Hydrophobic Surfaces: From Natural to Artificial. *Advanced Materials*, **14**, 57-60. <https://doi.org/10.1002/adma.200290020>
- [2] Cassie, A.B.D. and Baxter, S. (1944) Wettability of Porous Surfaces. *Transactions of the Faraday Society*, **40**, 546-551. <https://doi.org/10.1039/tf9444000546>
- [3] Wenzel, R.N. (1936) Resistance of Solid Surface to Wetting by Water. *Industrial & Engineering Chemistry*, **28**, 988-994. <https://doi.org/10.1021/ie50320a024>
- [4] Emami, B., Tafreshi, H.V., Gad-El-Hak, M., *et al.* (2012) Predicting Shape and Stability of Air-Water Interface on Superhydrophobic Surfaces Comprised of Pores with Arbitrary Shapes and Depths. *Applied Physics Letters*, **100**, 1-11. <https://doi.org/10.1063/1.3673619>
- [5] René, H., Ralf, H., Sebastian, A., *et al.* (2013) Wetting Resistance at Its Topographical Limit: The Benefit of Mushroom and Serif T Structures. *Langmuir: The ACS Journal of Surfaces & Colloids*, **29**, 1100-1112. <https://doi.org/10.1021/la304179b>
- [6] Liu, T. and Kim, C.J. (2014) Repellent Surfaces. Turning a Surface Super repellent even to Completely Wetting Liquids. *Science*, **346**, 1096-1100. <https://doi.org/10.1126/science.1254787>
- [7] Takahiro, K., Kenji, Y., Shigenori, F., *et al.* (2009) Coexistence and Transition between Cassie and Wenzel State on Pillared Hydrophobic Surface. *Proceedings of the National Academy of Sciences of the United States of America*, **106**, 8435-8440. <https://doi.org/10.1073/pnas.0902027106>
- [8] Jie, X. and Santanu, C. (2012) Design of Anti-Icing Coatings Using Supercooled Droplets as Nano-to-Microscale Probes. *Langmuir*, **28**, 4434-4446. <https://doi.org/10.1021/la2034565>
- [9] Solved with COMSOL Multiphysics 5.0, Filling of a Capillary Channel: Level Set and Phase Field Models.
- [10] Berthier, J., Loe-Mie, F., Tran, V., *et al.* (2009) On the Pinning of Interfaces on Micropillaredges. *Journal of Colloid and Interface Science*, **338**, 296-303. <https://doi.org/10.1016/j.jcis.2009.06.007>