

Hubble Diagram Test of 280 Supernovae Redshift Data

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ABSTRACT

We compare the Hubble diagram calculated from the observed redshift (RS)/magnitude (μ) data of 280 Supernovae in the RS range of $z = 0.0104$ to 8.1 with Hubble diagrams inferred on the basis of the exponential tired light and the Lambda Cold Dark Matter (Λ CDM) cosmological model. We show that the experimentally measured Hubble diagram follows clearly the exponential photon flight time (t_s)/RS relation, whilst the data calculated on the basis of the Λ CDM model exhibit poor agreement with the observed data.

KEYWORDS

Redshift/Magnitude Data Fitting; Linear Hubble Relation; Exponential Hubble Relation; Λ CDM Cosmological Model

1. Introduction

The basic assumption of the Lambda Cold Dark Matter (Λ CDM) cosmological model is that the universe is expanding, according to the Hubble's law [1], at a velocity of $v = zc = H_0 D_C$, where z is the redshift (RS), c is the velocity of light, H_0 is the Hubble constant, and D_C is the co-moving radial distance that can be derived from the observable z/μ data by (1).

$$D_C = \frac{10^{\frac{\mu+5}{5}}}{(z+1)} \times 3.085 \times 10^{18} \quad (1)$$

An important test of confidence in modeling the universal expansion is to compare the observed z/μ data with those derived on the basis of the Λ CDM model. The results presented in the literature, however, are not undisputed and are still a matter of debate. LaViolette [2] and more recently, López-Corredoira [3], Crawford [4], and Marosi [5-7] have shown that the static or slowly expanding universe models fit the observational data better than the data calculated on the basis of the presently prevailing Λ CDM model.

Such results, however, are usually refuted with the ar-

gument that the static universe contradicts many other cosmological observations, for example, the time dilation test and the cosmic microwave background (CMB) temperature versus RS test [8].

It is not the aim of this paper to argue in favor of or against either the expanding or static cosmological models. We only want to examine which of the two relations:

the linear Hubble's law or the exponential $e^{\frac{H_0 D_C}{c}} = z + 1$ fits the observational RS/ μ data more accurately.

We mean that the result of a proper data fitting procedure of reliable observational data cannot be ignored out of respect to the predictions of a theory. If facts contradict the theoretical expectations, then the only scientifically adequate answer can be that the underlying theory is at best, incomplete.

In this paper, we analyze the observed Hubble diagram compiled from 280 supernovae z/μ data in the range of $z = 0.0104$ to 8.1. We expect that in the high RS range, it should be possible to check more precisely whether the Hubble diagram follows a linear $z = H_0 D_C / c$ relation, or the exponential

$$e^{\frac{H_0 d}{c}} = e^{H_0 t_s} = z + 1 \quad (2)$$

relation; an effect that is perceptible only slightly in the $z < 1$ region.

2. Data Collection and Processing

In our analysis, we have included 171 gold-set data [9], 59 calibrated high-RS gamma-ray burst (GRB) data (Hymnium data set) and 50 low-RS GRBs obtained by Wei [10] from the 557 Union 2-compilation.

As the z/μ data are plagued by considerable scatter, similar to the procedure described in [5], the potential $\mu = a \times z^b$ function was used to perform a global fitting over the RS range of $z = 0.0104$ to 8.1.

As differences between the different cosmological models become more pronounced only in the linear t_S/z data representation, using Equations (2)-(4), the potential best fit data were converted into a t_S/z data set.

The photon flight time t_S was calculated from

$$t_s = \frac{D_c}{c} = \frac{10^5 \frac{\mu+5}{(z+1) \times 3 \times 10^{10}}}{c} \times 3.085 \times 10^{18} \quad (3)$$

In Equations (2) and (3), t_S means the flight time of the photons from the co-moving radial distance D_C to the observer, which should not be confused with the photon travel time (t) in an expanding universe. t_S means the flight time of photons between emission and reception, ($t_S = D_C/c$, c is the velocity of light), which is proportional to the D_C that is entered in the linear Hubble law.

The photon flight time t_S for the Λ CDM model was calculated with $H_0 = 72.6 \text{ km} \cdot \text{s}^{-1} \cdot \text{Mpc}^{-1}$, $\Omega_M = 0.266$, $\Omega_\Lambda = 0.732$ and $k = 0$ [11].

For the purpose of performing χ -squared tests in the high RS range of $t_S \times 10^{-14} = 6000$ to 11000 between the potential best fit and the t_S/z data calculated on the basis of the Λ CDM model we included 41 equidistant t_S/z data points in addition to the observed data.

The dimension of H_0 for the exponential function is expressed by the energy loss with time and it has the dimension $\text{Hz} \cdot \text{s}^{-1} \cdot \text{Hz}^{-1}$ instead of $\text{km} \cdot \text{s}^{-1} \cdot \text{Mpc}^{-1}$ as in the Λ CDM model.

Excel, Excel Solver and WinSTAT [12] software were used for the data fitting, refinement, and analysis and data presentation.

3. Results

The potential best fit curve of the 280 observed z/μ data points is shown in **Figure 1**.

Four outliers with standard deviation $> 3\sigma$ were identified in the z/μ data set and omitted from further regression analysis.

Results are shown in **Tables 1** and **2**.

It can be seen from **Table 2** that the omitted outliers have relatively little influence on the regression coefficient

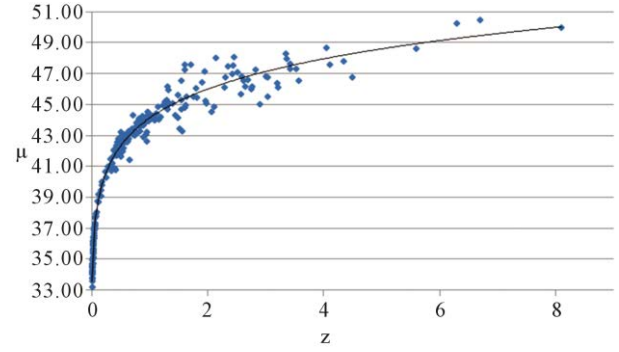


Figure 1. Solid line: potential $\mu = a \times z^b$ fit, diamonds: observed RS/ μ data.

Table 1. Outliers in the regression with 280 data points.

| Row | z-value | n*Sigma | P < 0.05 |
|-----|---------|---------|----------|
| 277 | 5.6 | 3.981 | 0.0152 |
| 278 | 6.29 | 4.573 | 0.0009 |
| 279 | 6.695 | 4.921 | 0.0001 |
| 280 | 8.1 | 6.127 | 6.218E-8 |

Table 2. Results of regression with $\mu = a \times z^b$ using 171, 276 and 280 z/μ data points.

| Data points | a | b | R2 | $\sum \chi_{\text{square}}$ | $P_{\chi_{\text{tests}}} \mu_{\text{obs}}/\mu_{\text{calc}}$ |
|--------------|-----------|----------|--------|-----------------------------|--|
| 171, Ref.[5] | 44.102 | 0.0593 | 0.9571 | 1.96634 | 1 |
| 276 | 44.109769 | 0.059883 | 0.9843 | 1.95407 | 1 |
| 280 | 44.1201 | 0.060005 | 0.9871 | 1.95407 | 1 |

coefficients a and b and that all the results for a and b lie within the very small error limits of $a \pm 0.02$ and $b \pm 0.0006$, respectively.

For further data treatment the potential best fit function obtained from 276 data points

$$\mu = 44.109769 \times z^{0.059883} \quad (4)$$

was used.

Tables 3-6 show the statistics of the fitting procedure with 276 data points.

4. The $t_S/(z + 1)$ Data Representation

Figure 2 shows the Hubble diagrams measured and calculated with $e^{2.024 \times 10^{-18} \times t_S} = z + 1$ in the range of $z + 1 = 1.0104$ to 5.35.

The goodness of fit indicators between the observed $t_S/(z + 1)$ data and the exponential $e^{2.024x}$ function for $z + 1 = 1.0104$ to 5.5, 6.5 and 9.1 are summarized in **Table 7**. The precise agreement between the measured and calculated data in the range of $z + 1 = 1.0104$ to 5.5 strongly supports the conclusion that the $t_S/(z + 1)$ function is exponential. It seems very likely that the small deviations at $z + 1 > 5.5$ are due to small systematic errors in distance measurements or to the calibration method at very high RSs.

Table 3. Descriptive statistics μ/z .

| | Valid cases | Mean | Std. error of mean | Variance | Std. Deviation | | |
|-------|-----------------------|-------------------------|--------------------|--------------|----------------|---------|-------|
| μ | 276 | 41.76894928 | 0.242237842 | 16.19545144 | 4.02435727 | | |
| | Variation coefficient | Rel. V. coefficient (%) | Skew | Kurtosis | Minimum | Maximum | Range |
| μ | 0.096348061 | 0.57994 | -0.622585385 | -0.735770108 | 33.21 | 48.68 | 15.47 |

Table 4. Descriptive statistics z/μ .

| | Valid cases | Mean | Std. error of mean | Variance | Std. Deviation | | |
|-----|-----------------------|-------------------------|--------------------|-------------|----------------|---------|--------|
| z | 276 | 0.880278623 | 0.056577708 | 0.883486217 | 0.939939475 | | |
| | Variation coefficient | Rel. V. coefficient (%) | Skew | Kurtosis | Minimum | Maximum | Range |
| z | 1.067774964 | 6.427249997 | 1.581919841 | 2.212059268 | 0.0104 | 4.5 | 4.4869 |

Table 5. Variable: μ ; grouped by z ; 95% confidence level.

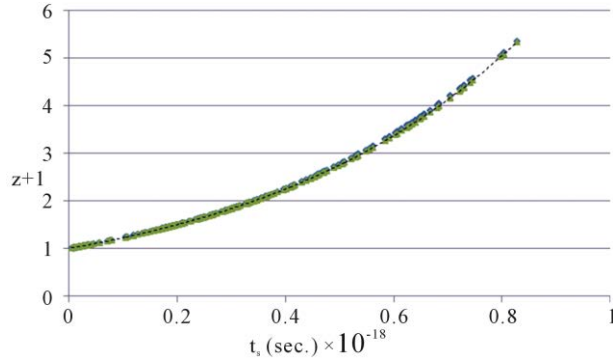
| μ | N | Mean | Conf. (\pm) | Std. Error | Std. Dev. |
|---------------|-----|-------------|-----------------|--------------|--------------|
| 33 to 34 | 4 | 0.012725 | 0.004668869 | 0.01467069 | 0.002934138 |
| 34 to 35 | 21 | 0.016414286 | 0.001003628 | 0.000481134 | 0.002204832 |
| 35 to 36 | 18 | 0.027077778 | 0.001658866 | 0.000786261 | 0.0003335823 |
| 36 to 37 | 15 | 0.41213333 | 0.003964023 | 0.001848213 | 0.007158099 |
| 37 to 38 | 10 | 0.0625 | 0.006904831 | 0.003052322 | 0.009652288 |
| 38 to 39 | 3 | 0.104666667 | 0.046628617 | 0.010837179 | 0.018770544 |
| 39 to 40 | 5 | 0.1612 | 0.027294023 | 0.009830565 | 0.021981811 |
| 40 to 41 | 6 | 0.330166667 | 0.077409426 | 0.030113581 | 0.073762908 |
| 41 to 42 | 16 | 0.42775 | 0.050113435 | 0.023511433 | 0.094045734 |
| 42 to 43 | 46 | 0.521336957 | 0.33007283 | 0.016388078 | 0.111149352 |
| 43 to 44 | 43 | 0.813906977 | 0.061365095 | 0.030407637 | 0.199396209 |
| 44 to 45 | 37 | 1.167456757 | 0.103807992 | 0.0511849998 | 0.311346189 |
| 45 to 46 | 15 | 1.823953333 | 0.321476572 | 0.149887432 | 0.580511527 |
| 46 to 47 | 19 | 2.640421053 | 0.355072924 | 0.169008137 | 0.73668939 |
| 47 to 48 | 14 | 2.799142857 | 0.529540171 | 0.245115646 | 0.917138766 |
| 48 to 49 | 4 | 2.99875 | 1.37869082 | 0.4332173 | 0.8664346 |
| Entire sample | 276 | 0.880278623 | 0.111380452 | 0.056577708 | 0.939939475 |

Table 6. Variable: z ; grouped by μ ; 95% confidence level.

| z | N | Mean | Conf. (\pm) | Std. Error | Std. Dev. |
|---------------|-----|------------|-----------------|--------------|-------------|
| 0.0 to 0.5 | 120 | 38.1235 | 0.578568683 | 0.292191652 | 3.200799178 |
| 0.5 to 1.0 | 76 | 3.29763158 | 0.148926165 | 0.074758293 | 0.65172769 |
| 1.0 to 1.5 | 28 | 44.6478571 | 0.233356837 | 0.113731045 | 0.601808122 |
| 1.5 to 2.0 | 18 | 45.6505555 | 0.62841576 | 0.297853401 | 1.263684958 |
| 2.0 to 2.5 | 9 | 46.6894444 | 1.000358877 | 0.433806188 | 1.301418564 |
| 2.5 to 3.0 | 10 | 46.331 | 0.485711017 | 0.214711434 | 0.678977172 |
| 3.0 to 3.5 | 9 | 46.975 | 0.693006582 | 0.300522693 | 0.901568078 |
| 3.5 to 4.0 | 2 | 46.93 | 4.828358221 | 0.38 | 0.537401154 |
| 4.0 to 4.5 | 3 | 48.0166666 | 1.45068823 | 0.3371611353 | 0.583980593 |
| 4.5 to 5.0 | 1 | 46.74 | - | - | - |
| Entire sample | 276 | 41.7689492 | 0.476876166 | 0.242237842 | 4.02435727 |

Table 7. Goodness of fit indicators.

| Data points | $z + 1$ range | R^2 | Std. error | Std. dev. | $\sum \chi^2$ | P |
|-------------|---------------|---------|------------|-----------|---------------|---|
| 276 | 1.0104 - 5.5 | 0.99996 | 0.006190 | 0.933469 | 0.007723956 | 1 |
| 278 | 1.0104 - 6.5 | 0.99985 | 0.019795 | 1.027633 | 0.010656491 | 1 |
| 280 | 1.0104 - 9.1 | 0.99838 | 0.046957 | 1.173129 | 0.088924843 | 1 |

**Figure 2.** Observed data (diamonds), data calculated with $e^{2.024 \times 10^{-18} t_s}$ (triangles), trendline with $e^{2.024x}$ (dashed line).

t_s/z Diagram in the Range of $z = 0.0104 - 8.1$, Comparison with the Λ CDM Model

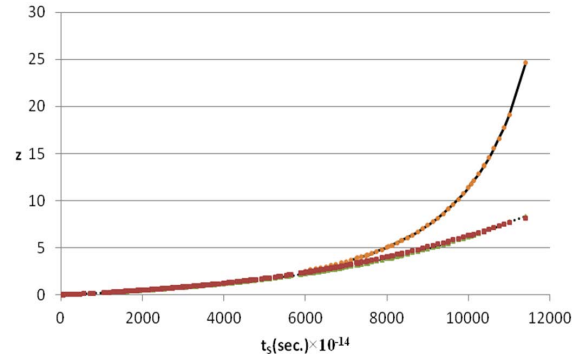
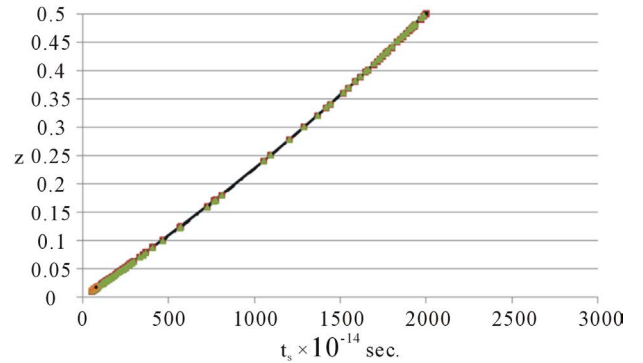
Figure 3 shows the t_s/z diagram in the range of $z = 0.0104$ to 8.1 calculated using Equation (4) with the observed z/μ data set (squares), the exponential function $z = e^{2.024 \times 10^{-18} t_s} - 1$ (triangles), and the t_s/z relation derived from the Λ CDM model (circles) with $H_0 = 72.6 \text{ km s}^{-1} \text{ Mpc}^{-1}$, $\Omega_M = 0.266$, $\Omega_\Lambda = 0.732$ and $k = 0$.

One can see from **Figure 3** that, similar to the plot shown in **Figure 2**, the curves calculated from the best fit and the exponential function $z = e^{2.024 \times 10^{-18} t_s} - 1$ are nearly concurrent over the entire range of z , Pchi square = 1, whilst at $z > 2$ the t_s/z data calculated based on the Λ CDM model show clearly a different slope and depart considerably from both, the linear and the exponential functions. The χ -square test indicates statistical significance between the observed t_s/μ and the calculated Λ CDM data of $P = 0.0173$, indicating that from a statistical point of view, the two models are essentially different.

At RSs $z < 0.3$ (**Figure 4**), the t_s/z curves for the potential best fit, the exponential function, and the Λ CDM model can be fitted with the linear function $z = 0.000228725 \times t_s - 0.00332331$ ($R^2 = 0.9989$) with good approximation. The linear approximation, however, is deceiving. As can be seen in **Figure 3**, that at high RSs, the best-fit and the exponential curves follow strictly the exponential energy depletion relationship.

5. Conclusions

The most impressive result of the Hubble diagram test is

**Figure 3.** Redshift of type Ia supernovae as a function of $t_s = D_c/c$. Squares: t_s/z data inferred from the potential best-fit curve of the observed z/μ diagram. Triangles: the exponential t_s/z relation with $H_0 = 2.024 \times 10^{-18}$. Circles: t_s/z relation derived from the Λ CDM model with $H_0 = 72.6 \text{ km} \cdot \text{s}^{-1} \cdot \text{Mpc}^{-1}$.**Figure 4.** The “linear” t_s/z relation in the low RS region for the potential best fit and the exponential function, and for the t_s/z data calculated based on the Λ CDM model.

that the t_s/z relation obtained from the potential best fit data can be expressed nearly exactly by the exponential formula $e^{2.024 \times 10^{-18} t_s} = z + 1$ over the entire range of $z = 0.01$ to 8.1 .

In contrast, in the RS range $z > 2$ the t_s/z curve derived from the Λ CDM model with $H_0 = 72.6 \text{ km} \cdot \text{s}^{-1} \cdot \text{Mpc}^{-1}$, $\Omega_M = 0.266$, $\Omega_\Lambda = 0.732$ and $k = 0$, shows poor agreement with the observed data. The χ -square test indicates statistical significance between the observational potential fit and the calculated Λ CDM data of $P = 0.0173$, indicating that from a statistical point of view the two models are essentially different.

Based on the results presented in this paper, a reconsideration of the Λ CDM model appears warranted.

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