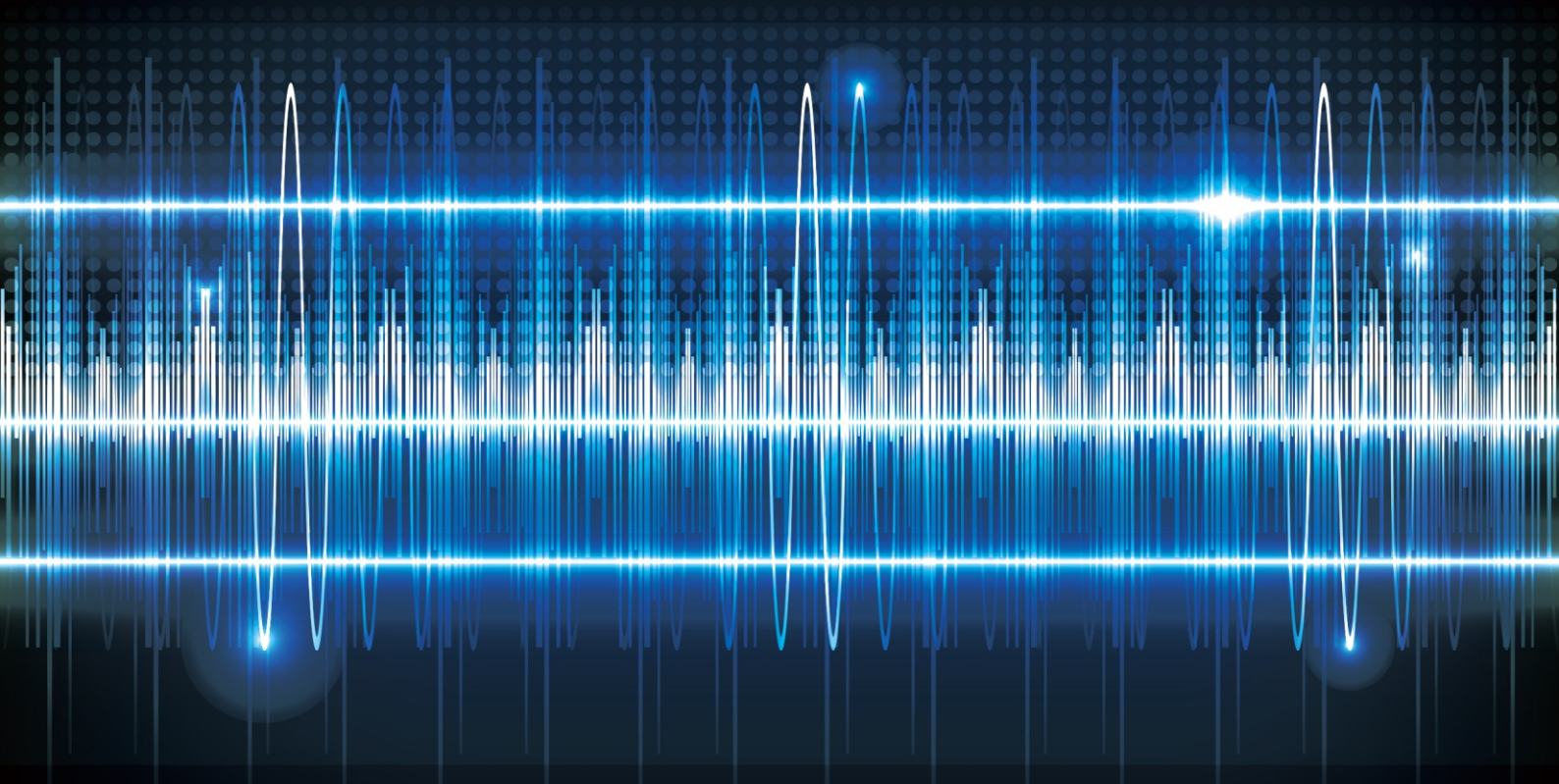


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Analysis of Certain Fruit Juices Using Microwave Techniques

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Abstract

This paper discusses a novel method of analyzing the dielectric properties of juices at various frequencies of microwave band. The measurements were performed using a rectangular cavity with the perturbation technique at the S-band of microwave frequency with fresh juices made instantly from the fruits as well as collected from the packed fruit juices available in the market with the normal added preservatives in order to keep the increase of the shelf life. From the results, it is observed in the certain dielectric properties, the fresh juices samples and packed juices samples were varying and also varying over a period of time. This measurement method is simple and quick and can be used over a range of juices. These results prove a new method of determining the quality control of juices using microwave principles.

Keywords

Fruit, Juice, Cavity Perturbation, Shelf Life, Quality Control

1. Introduction

Fruits are normally sweet-tasting plant product, sometimes with seeds or without seeds, and most of them are edible in the raw state. Humans and animals make use of fruits as a source of food. In terms of the nutrition value, a fruit basically consists of macronutrients and micronutrients. The macronutrients like fiber and micronutrients like minerals, vitamins C, thiamin, riboflavin, B-6, niacin, folate, A, and E. Phytochemicals like polyphenolics, carotenoids, and glucosinolates, will add nutritional value to a fruit [1] [2]. Mechanically squeezing or macerating a fruit will lead the preparation of a juice. Also though juice will have reduced macronutritional value, it can be made available at all times by

addition of preservatives while fruits are available only at the designated period of the crop. So packed fruit juice have very high commercial value in today's world. But packed fruit juice will have added preservatives like ascorbic acid, carbon dioxide and in some cases colorants in order to increase its shelf life. The present study is an attempt to look into the dielectric properties of fresh juice as well as packed fruit juice at microwave frequencies.

At electromagnetic field, the dielectric property depends on water content, frequency, intensity of the field as well as temperature. As the temperature increases, more and more dipoles are free to oscillate, thereby there will be an increase in dielectric characteristics. The dielectric property is represented by dielectric constant and dielectric loss factor. The dielectric constant decides how much incident energy can go through the material and dielectric loss determines how much energy can be absorbed within the material. In reality, the energy absorption as well as heat dissipation in a material is purely on the basis of specific characteristics like density, water content, composition and surface structure as well as bulk characteristics like shape, size and porosity [3].

The microwave technology is used in various applications in food industry because of the energy is directly transferred to the materials which have moisture content rather than the conventional conduction mechanism which is a slow process. Over the conventional processing methods, microwave technology provides a rapid, clean and efficient heating over a wide range of temperatures. Electromagnetic energy is directly transferred to and absorbed by the moisture content material being processed in microwave processing. Thus, microwave energy can penetrate, thereby producing heat makes microwave technology better advantages when compared with other conventional heat processing techniques so that it is used in reheating, pasteurization, curing, freeze-drying, thawing, sterilization, tempering, baking, pre-cooking, dehydration and blanching [4] [5] [6] [7] [8].

Various measurement methods are used to measure the complex permittivity of a material and the method of choice depends on factors like the nature of the samples under study and the frequency range used [9] [10] [11] [12]. When only very small volumes of the sample are available, the cavity perturbation technique is an attractive option as it requires only minute volumes for the measurement [13]. This makes it suitable for the dielectric study of juices needed. However, no data is currently available for the complex permittivity of fresh juices and packed juices in the literature. In this study, the rectangular cavity perturbation technique has been employed for the measurement of the dielectric properties of freshly prepared juices as well as packed juices available in the market with the chemically added preservatives in order to increase its shelf life. The frequency of measurement is made between 2 to 4 GHz. It is observed that a more appreciable change in the dielectric properties of packed fruit juice than the fresh samples at the extended period of time. This microwave measurement technique is very simple, quick and the results are promising as this new method can be

used in the quality control of packed fruit juices.

2. Sample Preparation

The study was conducted on pineapple, orange, naartjie, grape, apple, lemon, guava, lychee, strawberry and mango. Fresh fruits are obtained from the market and made into juice by mechanically squeezing just before the measurements. Similarly the packed fruit juice from the market which are kept at temperature at 4°C to 8°C were used for the measurements. The juices were filled into the capillary bulb for microwave measurements and the measurements were made at 20°C.

3. Material and Methods

The experimental set-up consists of a transmission type S-band rectangular cavity resonator, Rohde & Schwarz ZVL network analyzer. The cavity resonator is a transmission line with one or both ends closed. **Figure 1** shows the experimental setup. The numbers of resonant frequencies are determined by the length of the resonator. The resonator in this set-up is excited in the TE_{10p} mode. The sample holder which is made of glass in the form of a capillary tube flared to a disk shaped bulb at the bottom is placed into the cavity through the non-radiating cavity slot, at broader side of the cavity which can facilitate the easy movement of the holder. The resonant frequency f_0 and the corresponding quality factor Q_0 of the cavity at each resonant peak with the empty sample holder placed at the maximum electric field are noted. The same holder filled with known amount of sample under study is again introduced into the cavity resonator through the non-radiating slot. The resonant frequencies of the sample loaded cavity are selected and the position of the sample is adjusted for maximum perturbation (*i.e.* maximum shift of resonant frequency with minimum amplitude for the peak). The new resonant frequency f_s and the quality factor Q_s are noted. The same procedure is repeated for other resonant frequencies. The measurement is made from 2 GHz to 4 GHz of S-band in microwave region [14].

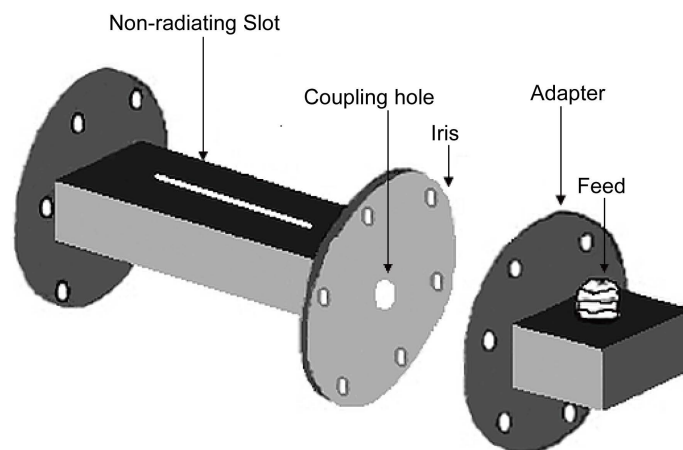


Figure 1. Experimental set up.

4. Theory

Theoretical analysis for the determination of complex permittivity of materials using rectangular waveguide cavity: When a small sample is inserted in a cavity which has the electric field E_0 and magnetic field H_0 in the unperturbed state, the fields in the interior of the object are E and H . Beginning with Maxwell's equations, Bethe and Schwinger [15] obtained an expression for the resonant frequency shift. For lossless sample, the variation of resonant frequency is given by Harrington [16] as

$$\frac{\omega - \omega_0}{\omega} = - \frac{\int (\Delta \varepsilon E \cdot E_0^* + \Delta \mu H \cdot H_0^*) d\tau}{\int (\varepsilon E \cdot E_0^* + \mu H \cdot H_0^*) d\tau} \quad (1)$$

where ε and μ are the permittivity and permeability of the medium in the unperturbed cavity respectively and $d\tau$ is the elemental volume. $\Delta \varepsilon$ and $\Delta \mu$ are changes in the above quantities due to introduction of the sample in the cavity. Without affecting the generality of Maxwell's equations, the complex frequency shift due to a lossy sample in the cavity is given by Waldron [17]

$$-\frac{\delta \Omega}{\Omega} \approx \frac{(\bar{\varepsilon}_r - 1) \varepsilon_0 \int_{V_s} E \cdot E_0^* dV + (\bar{\mu}_r - 1) \mu_0 \int_{V_s} H \cdot H_0^* dV}{\int_{V_c} (D_0 \cdot E_0^* + B_0 \cdot H_0^*) dV} \quad (2)$$

Two approximations are made in applying Equation (2), based on the assumption that fields in the empty part of the cavity are negligibly changed by the insertion of the sample and that the fields in the sample are uniform over its volume. Both these assumptions can be considered valid if the object is sufficiently small relative to the resonant wavelength. The negative sign in Equation (2) indicates that by introducing the sample the resonant frequency is lowered. Because the permittivity of practical materials is complex, the resonant frequency should also be considered as complex. In Equation (2), the $\delta \Omega$ is the complex frequency shift. B_0 , H_0 , D_0 and E_0 are the fields in the unperturbed cavity. E and H are the fields in the interior of the sample. $\bar{\varepsilon}_r = \bar{\varepsilon}'_r - j\bar{\varepsilon}''_r$ and $\bar{\mu}_r = \bar{\mu}'_r - j\bar{\mu}''_r$ and V_c and V_s are the volumes of the cavity and sample respectively. In terms of energy, the numerator of Equation (2) represents the energy stored in the sample and the denominator represents the total energy stored in the cavity. The total energy $W = W_e + W_m = 2W_e = 2W_m$. When a dielectric sample is introduced at the position of maximum electric field only the first term in the numerator is significant, since a small change in ε at a point of zero electric field or a small change in μ at a point of zero magnetic field does not change the resonance frequency. Thus Equation (2) can be reduced to

$$-\frac{\delta \Omega}{\Omega} \approx \frac{(\bar{\varepsilon}_r - 1) \int_{V_s} E \cdot E_{0\max}^* dV}{2 \int_{V_c} |E_0|^2 dV} \quad (3)$$

Let Q_0 be the quality factor of the cavity in the unperturbed condition and

Q_s the Q-factor of the cavity loaded with the object. The complex frequency shift is related to measurable quantities by [18]

$$\frac{\delta\Omega}{\Omega} \approx \frac{\delta\omega}{\omega} + \frac{j}{2} \left[\frac{1}{Q_s} - \frac{1}{Q_0} \right] \quad (4)$$

Equating the real and imaginary terms of Equations (3) and (4) we get

$$-\frac{f_s - f_0}{f_s} = \frac{(\varepsilon_r' - 1) \int_{V_s} E \cdot E_{0\max}^* dV}{2 \int_{V_c} |E_0|^2 dV} \quad (5)$$

$$\frac{1}{2} \left[\frac{1}{Q_s} - \frac{1}{Q_0} \right] = \frac{\varepsilon_r'' \int_{V_s} E \cdot E_{0\max}^* dV}{2 \int_{V_c} |E_0|^2 dV} \quad (6)$$

We may assume that $E \approx E_0$ and the value of E_0 in TE_{10p} mode as $E_0 = E_{0\max} \sin\left(\frac{m\pi x}{a}\right) \sin\left(\frac{p\pi z}{d}\right)$ where a is the broader dimension of the waveguide and d is the length of the cavity. Integrating and rearranging the above equations we get

$$\varepsilon_r' - 1 = \frac{f_0 - f_s}{2f_s} \left(\frac{V_c}{V_s} \right) \quad (7)$$

$$\varepsilon_r'' = \frac{V_c}{4V_s} \left[\frac{1}{Q_s} - \frac{1}{Q_0} \right] \quad (8)$$

If the frequency shift is measured from the resonance frequency f_t of the cavity loaded with empty capillary rather than that with empty cavity alone the above equations become

$$\varepsilon_r' - 1 = \frac{f_t - f_s}{2f_s} \left(\frac{V_c}{V_s} \right) \quad (9)$$

$$\varepsilon_r'' = \frac{V_c}{4V_s} \left[\frac{1}{Q_s} - \frac{1}{Q_t} \right] \quad (10)$$

Q_t is the quality factor of the cavity loaded with empty capillary. f_s and Q_s are the resonance frequency and quality factor of the cavity loaded with capillary containing the sample material.

Theory for the determination of conductivity of the materials:

For a dielectric material having non-zero conductivity, the Ampere's law in phasor form as

$$\nabla \times H = (\sigma + j\omega\bar{\varepsilon})E = (\sigma + j\omega\varepsilon'')E + j\omega\varepsilon'E \quad (11)$$

where $\varepsilon = \varepsilon' - j\varepsilon''$ is the absolute permittivity of the medium.

The loss tangent

$$\tan \delta = \frac{\sigma + \omega \epsilon''}{\omega \epsilon'} \tag{12a}$$

For dielectrics $\sigma = 0$ then Equation (12a) becomes

$$\tan \delta = \frac{\epsilon''}{\epsilon'} \tag{12b}$$

$\sigma_e = \sigma + \omega \epsilon''$, is the effective conductivity of the medium.

But

$$\tan \delta = \frac{1}{Q_m} = \left[\frac{1}{Q_s} - \frac{1}{Q_i} \right] \tag{13}$$

where is Q_m is the loaded Q -factor of the cavity with sample alone.

The effective conductivity

$$\sigma_e = \frac{\omega \epsilon'}{Q_m} = \frac{\omega \epsilon_r' \epsilon_0}{Q_m} \tag{14}$$

When σ is very small, the effective conductivity is reduced to

$$\sigma_e = \omega \epsilon'' = 2\pi f \epsilon_0 \epsilon_r'' \tag{15}$$

5. Results and Discussion

The microwave experiment in fruit juices were done using cavity perturbation technique collected from freshly prepared juices as well as packed juices and the results were shown in Figures 2-13. From Figures 2-4 it is observed that freshly prepared juices as well as packed juices exhibit almost similar ranges of dielectric constant even though the measurements were made at different intervals of time. From Figures 5-7 shows the dielectric loss or conductivity of freshly prepared juices. These results are shows that the conductivities are consistent over the extended period of time. From Figures 8-10 shows the dielectric constant of packed

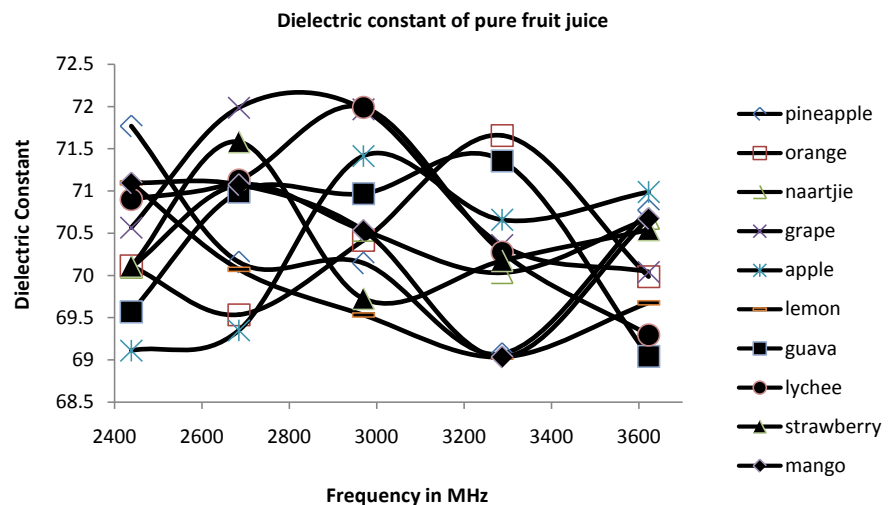


Figure 2. Variation of dielectric constant of pure fruit juice.

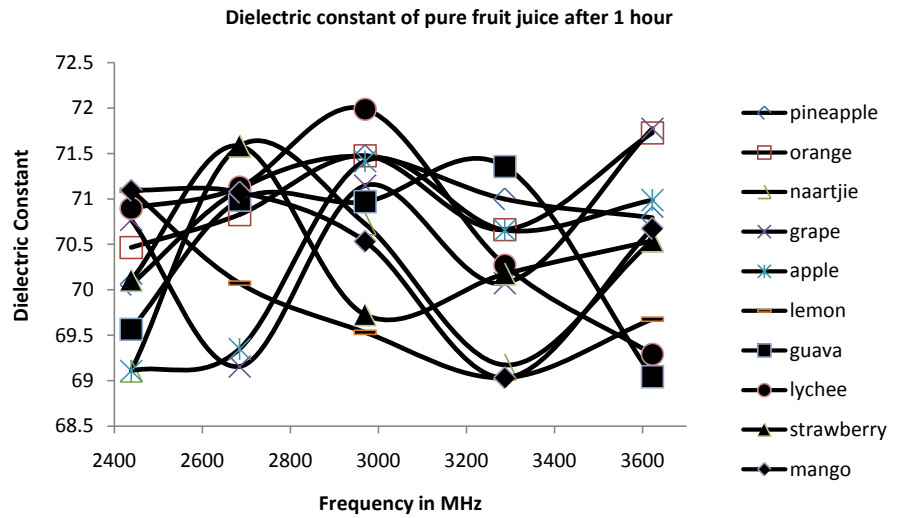


Figure 3. Variation of dielectric constant of pure fruit juice after 1 hour.

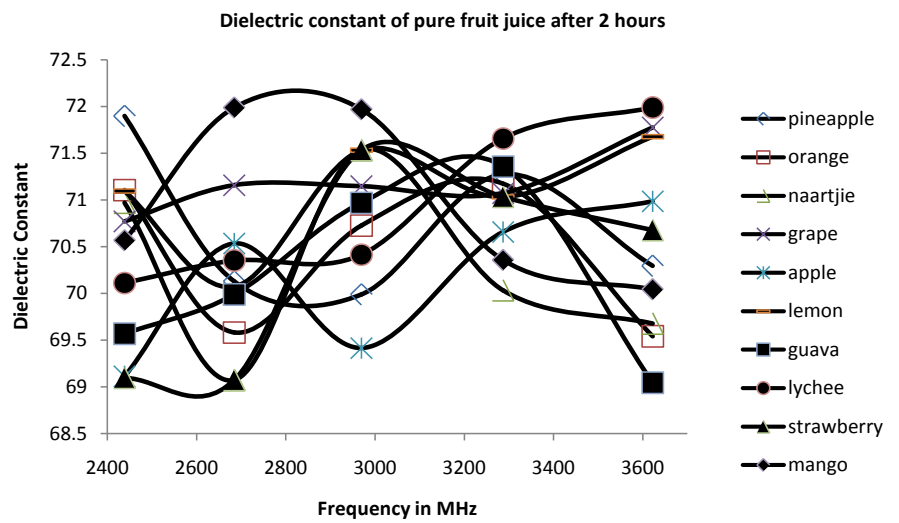


Figure 4. Variation of dielectric constant of pure fruit juice after 2 hours.

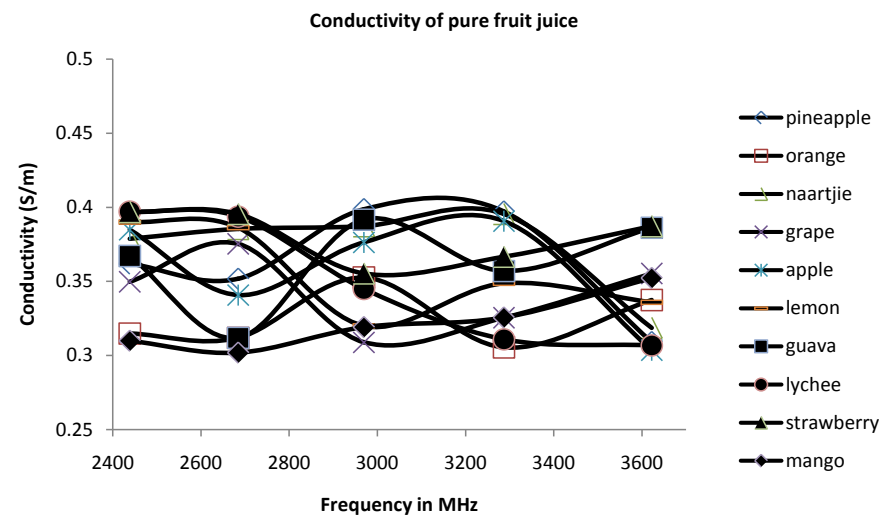


Figure 5. Variation of conductivity of pure fruit juice.

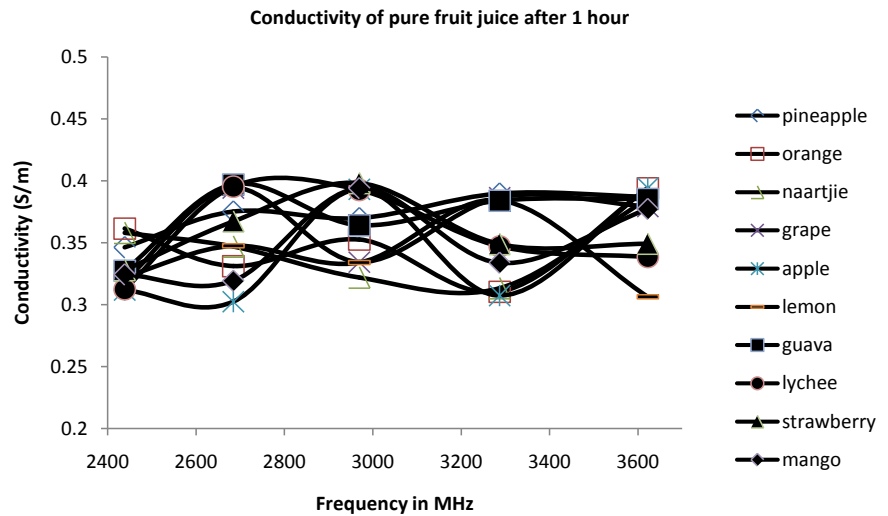


Figure 6. Variation of conductivity of pure fruit juice after 1 hour

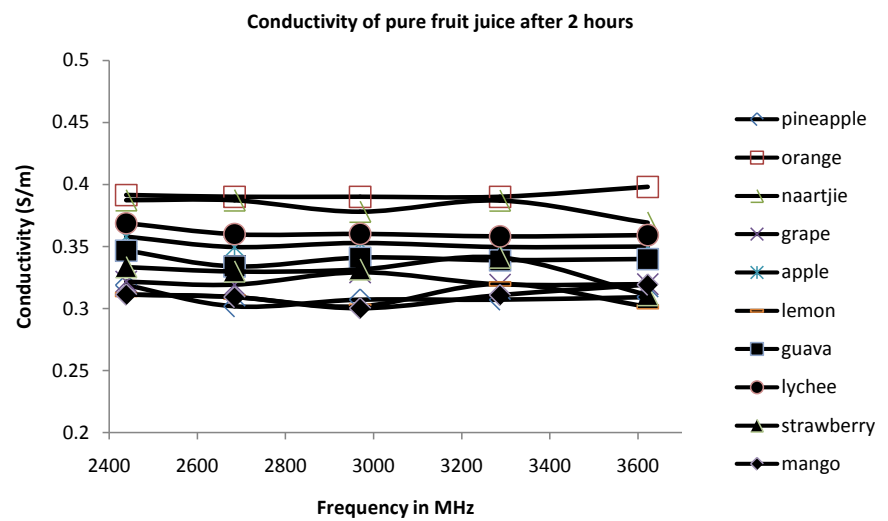


Figure 7. Variation of conductivity of pure fruit juice after 2 hours.

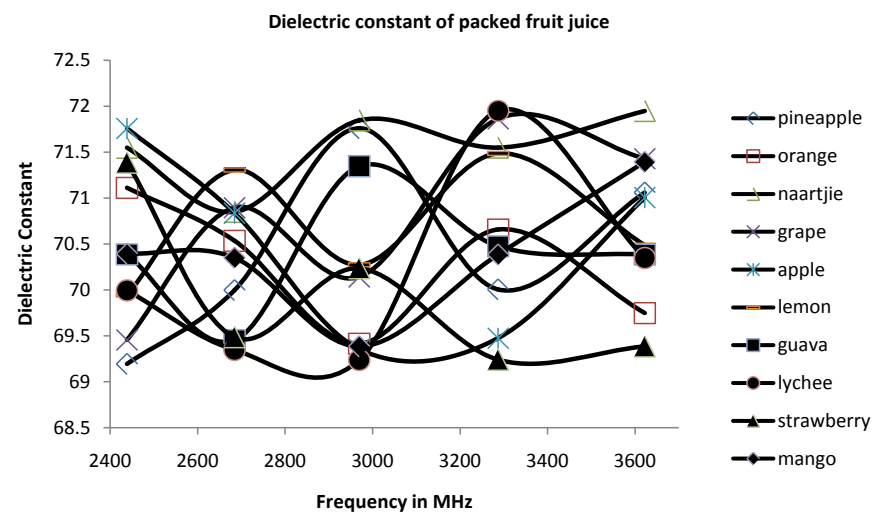


Figure 8. Variation of dielectric constant of packed fruit juice.

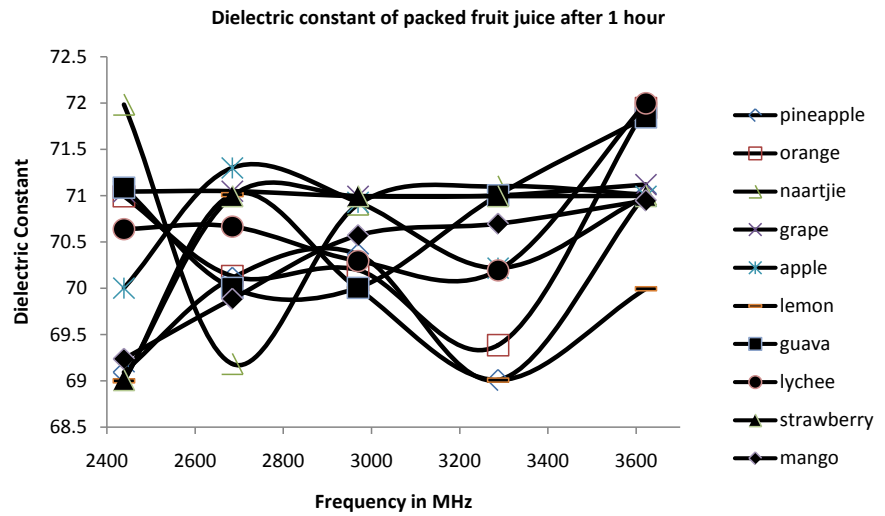


Figure 9. Variation of dielectric constant of packed fruit juice after 1 hour.

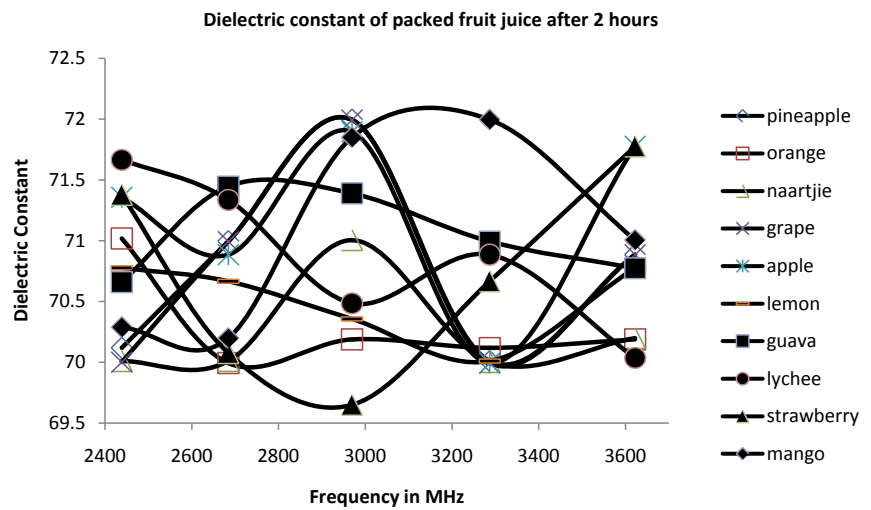


Figure 10. Variation of dielectric constant of packed fruit juice after 2 hours.

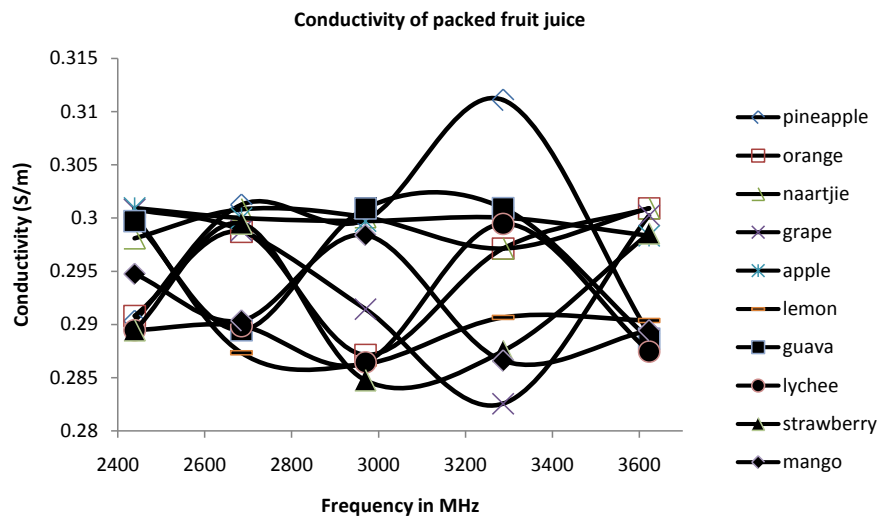


Figure 11. Variation of conductivity of packed fruit juice.

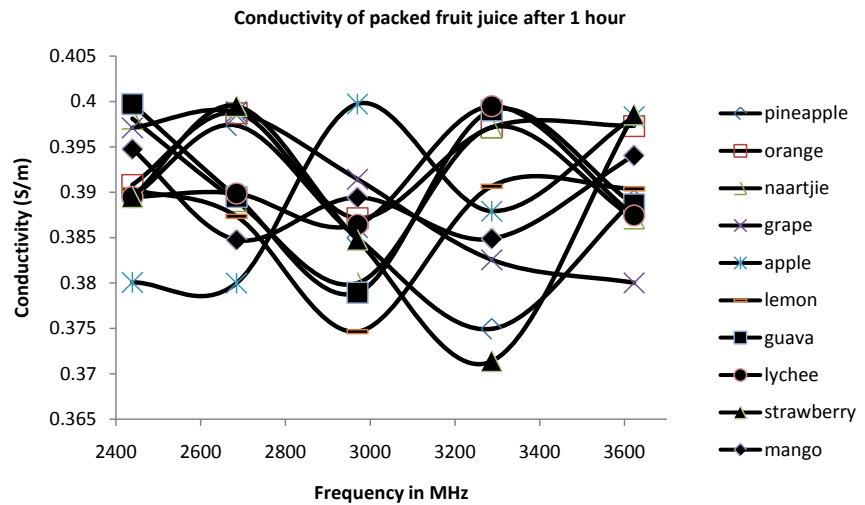


Figure 12. Variation of conductivity of packed fruit juice after 1 hour.

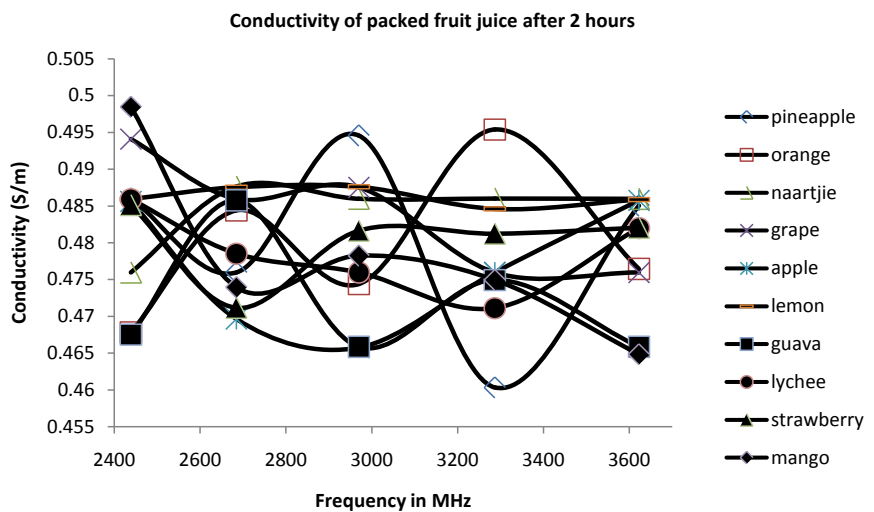


Figure 13. Variation of conductivity of packed fruit juice after 2 hours.

fruit juice at different time intervals. These results suggest that the dielectric constant is similar for the different intervals of time. From **Figures 11-13** shows the conductivities of packed fruit juice over the extended period of time. These results show that there is a distinct variation of conductivities over the period of time as well as with that of freshly prepared juices. Packed juices exhibit higher conductivity than the freshly prepared juices because of the presence of added preservatives which is necessary to increase its shelf life. So it is very critical that packed fruit juice should be consumed immediately after it is opened. Thus in the S band of microwave (ISM band), freshly prepared fruit juices and packed fruit juice were studied and exhibit distinct variation of conductivity even after certain intervals of time.

6. Conclusion

The microwave characterization has been performed in the freshly prepared

juices as well as packed fruit juices using the cavity perturbation technique. The cavity perturbation technique is quick, simple, and accurate and it requires very low volume of sample for measuring the dielectric properties of samples like juices. The limitation of this method is that measurement can lead to more accurate results on liquid samples than on solid samples. From the results, it is observed in the certain dielectric properties, the fresh juices samples and packed juices samples were varying and also varying over a period of time. This measurement method is simple and quick and can be used over a range of juices. These results prove a new method of determining the quality control of juices using microwave principles.

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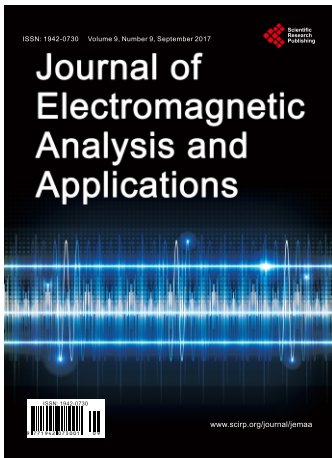
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- Electrical Machine
- Electrocardiograph (ECG)
- Electroencephalograph (EEG)
- Electromagnetic Breaker
- Electromagnetic Compatibility (EMC)
- Electromagnetic Compatibility and Electromagnetic Environment
- Electromagnetic Devices
- Electromagnetic Field Theory
- Electromagnetic Interferences (EMI)
- Electromagnetic Inverse Problems
- Electromagnetic Launch
- Electromagnetic Material Modelling
- Electromagnetic Measurement Technology and Instruments
- Electromagnetic Nondestructive Testing
- Electromagnetic Numerical Analysis
- Electromagnetic Physics
- Electromagnetic Solid Mechanics
- Electromagnetic Structure Optimization
- Electromagnetism and Biological Tissues
- Electromagnetism and Medical Devices
- Electromagnetism in Medical Applications

- Electromyography
- Environmental Pollution by Electromagnetics
- Fiber Optics
- High Frequency Techniques
- Integrated Optics
- Magnetic Fluid
- Measuring Technique of Radiated Electromagnetic
- Medical Applications
- Moving Conductor Eddy Current Problems
- Multiphysics Coupled Problems
- Noise and Signals
- Noise Reduction
- Optical and Millimeter Wave Techniques
- Printed Circuits
- Quasi-Static Fields
- Radar Measurements and Applications
- Radiated Electromagnetic
- Scattering and Diffraction
- Serges (Thunderbolts or Static Electricity)
- Solid State Devices and Circuits
- Static Fields
- Static Magnetic
- Test Electromagnetic Analysis Method (Team) Workshop Benchmark Problems
- The Problems of the Propagation of Electromagnetic Waves
- Waveguides

Notes for Intending Authors

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