

# The Color of the Sky

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## ABSTRACT

The color of the sky in day-time and at twilight is studied by means of spectroscopy, which provides an unambiguous way to understand and quantify why a sky is blue, pink, or red. The colors a daylight sky can take primarily owe to Rayleigh extinction and ozone absorption. Spectra of the sky illuminated by the sun can generally be represented by a generic analytical expression which involves the Rayleigh function  $R \propto \frac{1}{\lambda^4} e^{-\frac{a}{\lambda^4}}$ , ozone absorption, and, to a lesser extend, aerosol extinction. This study is based on a representative sample of spectra selected from a few hundred observations taken in different places, times, and dates, with a portable fiber spectrometer.

**Keywords:** Sky Color; Rayleigh Extinction; Ozone; Aerosols; Radiative Transfer; Spectral Analysis

## 1. Introduction

Two major causes of the extinction of light by the atmosphere, extinction by air molecules (nitrogen mainly) and ozone absorption, were discovered in the second half of the XIX<sup>th</sup> century when scientists<sup>1</sup> questioned why the sky is blue. A third important cause of atmospheric extinction, extinction by aerosols, has been observed for centuries to affect sky color, for instance the reddish hazes after a volcanic eruption or sand-winds [7,8]. Aerosol science also developed in the XIX<sup>th</sup> century although the theoretical understanding of how these particles extinguish light had to wait for Mie theory in 1908 [9]. Today, information about the atmosphere is primarily derived from spectral analysis. Occultation spectra of stars or of the sun are inverted to give the exact column densities of nitrogen, ozone, and aerosols, along the path of light (**Figure 1**).

Surprisingly, no equivalent method has been used for scattered sunlight. Despite a large number of investigations on the color of the sky<sup>2</sup>, little effort has been made to relate a color to its spectrum in order to understand why is a sky red, the difference between day-time and twilight blues, or the reason for a green-flash. Since about sixty years investigations on the color of the sky are based on models of the atmosphere [10]. These models

agree that the blue color of the sky in day-time is due to Rayleigh scattering and that blue nuances at twilight, when the sun is low on the horizon, require an additional ozone absorption [11-16]. A conclusion that Wulf, Moore, and Melvin [17] had already reached in 1934 from what may be considered as the very first published spectral observation of a twilight sky. Recently, Lee, Meyer, and Hoeppe [10] have also addressed the question of the influence of ozone on the blue of the sky from a calorimetric analysis of digital images of the sky.

At the other end of the visible spectrum the reason for the red color of the horizon at sunrise or sunset is still an open question. Rayleigh extinction is often thought to explain red skies<sup>3</sup>, but models find that an additional contribution of either multiple scattering or aerosol extinction is needed [13,16].

A first attempt to analyze the spectrum of a red horizon was made in 2001-2003 by M. Fujii and myself [19]. The spectra observed by M. Fujii had exactly the shape of occultation spectra of the sun (which I knew from [20]) multiplied by  $1/\lambda^4$ . They should simply be interpreted as sunlight extinguished by the atmosphere Rayleigh-scattered by air-molecules in the direction of the observer. The spectra were slightly (a few degrees) above the red-

<sup>3</sup>“Scattering decreases, in fact, as  $1/\lambda^4$ . Blue and violet light are thus scattered much more than red and orange, which is why the sky looks blue. At sunset, the Sun’s rays pass through a maximum length of atmosphere. Much of the blue has been taken out by the scattering. The light that reaches the surface of the Earth, and reflects off clouds and haze, is thus lacking in blue which is why sunsets appear reddish.” ([18], p. 945). See also: <http://www.sciencedaily.com/releases/2007/11/071108135522.html>

<sup>1</sup>Physicists and chemists, starting with Clausius and Rayleigh ([1,2] and references therein) followed by Lallemand, Hartley, Chappuis [3-6].

<sup>2</sup>See for instance the 2005 issue of Applied Optics, *n*<sup>o</sup> 44 (No. 27, Sept. 20), and references therein.

dest part of the horizon, but the article nevertheless suggests that ozone, which absorbs the orange-green part of the spectrum, plays a role in the red color of the sky as central as for blue twilights. The article also proved that contrary to widespread opinion observed spectra of scattered sunlight can indeed be analyzed and quantified in the same way as extinguished sunlight is, independently from a modeling of the atmosphere. The obvious advantage is to free the conclusions from arbitrary inputs, either parameters such as column densities of nitrogen or ozone along the path of light, which are precisely the quantities observation should retrieve, or the specific geometry of the scattering and the structure of the atmosphere, which are always difficult to ascertain.

Few such studies exist. Twilight spectra observed at the Mount Paranal ESO Observatory (Chile) by Patat *et al.* [21] highlight the importance of ozone, but the authors made no attempt to retrieve quantitative information from their spectra. A more consistent analysis of sky spectra observed with an Ocean Optics portable spectrometer was recently published by D. Fosbury, G. Koch, and J. Koch [22].

Occultation observations of stars (**Figure 1**), or the early blue sky observations by Wulf, Moore, and Melvin [17], show how powerful a tool spectral observations can be in the analysis of sunlight (extinguished or scattered) viewed from the earth. The present study furthers the work initiated in [19,20] to characterize extinguished and scattered sunlight from their spectrum. The observations are described in Section 2. Basic properties of extinguished sunlight are recalled in Section 3. Sections 4-6 analyze the properties of spectra of blue and red skies. These results are discussed in Section 7.

## 2. Observations

The spectrometer used is a miniature Ocean Optics USB2000 VIS-NIR portable fiber spectrometer. On several occasions between June 2010 and March 2011, spectra were recorded on a portable computer. Sunsets, sunrises, blue daytime and twilight skies, were repeatedly observed from different locations in New England (USA). For sunsets and sunrises I focused on the variations with time of the spectrum, which provided the reasons of the change of color in the sky and the best fit to apply to the spectra. The consistency of the red spectra was checked against sunrise spectra (with a 1.2" spatial resolution) observed with the 1.5 m telescope of the Fred Whipple Observatory (USA). Blue skies were observed at several distances from the sun and different elevations. The spectra selected for the article are representative of different sky colors and offer a consistent picture of the reasons why a sky turns from blue to red.

The spectral resolution is 2 Å. The instrument sensi-

tivity is limited compared to more sophisticated models but proved to be sufficient for the sky observations, as long as they were not too bright (no direct light from the sun) and not too dimmed (no night sky). The fiber's aperture is given to be 25°; in practice the fiber is reactive to direction of a source of light which contrasts with its surrounding, which is the case at sunrise or sunset. Exposure time varied from 0.1 s (blue skies) to 60 s (first twilit spectra at sunrise).

Plots are limited to the 400 - 900 nm (1.1 - 2.5  $\mu\text{m}^{-1}$ ) region. The spectra are presented on an arbitrary scale set to be 1 at 2.5  $\mu\text{m}^{-1}$  (400 nm) for the blue zenith sky spectrum of **Figure 2**.

## 3. Extinction of Direct Sunlight

At visible wavelengths a beam of light crossing a clear (no cloud) atmosphere experiences three major sources of extinction (**Figure 1**): Rayleigh and aerosol extinctions and ozone absorption. The spectrum  $S_0$  of sunlight observed through a layer of atmosphere is attenuated by a factor

$$\frac{S}{S_0} \propto T = e^{-\frac{a}{\lambda^4}} e^{-\frac{b}{\lambda^p}} e^{-N_{O_3}\sigma_{O_3}} \quad (1)$$

and the spectrum of the sun in any location in the atmosphere must be  $\propto TS_0$ . In Equation (1) the exponentials stand for Rayleigh extinction by nitrogen, aerosol extinction, and ozone absorption.  $N_{O_3}$  is the ozone column density along the path of sunlight,  $\sigma_{O_3}$  is the wavelength dependent ozone absorption cross-section.

Aerosols provide the  $e^{-\frac{b}{\lambda^p}}$  term of Equation (1),  $b$  being proportionate to the column density of aerosols along the path of sunlight, from outside the atmosphere to the point of observation. Aerosols scatter and extinguish light as  $1/\lambda^p$ , with  $p$  close to 1 ( $p \approx 1.3$ , [24,25]); scattering by aerosols is strongly oriented in the forward direction.

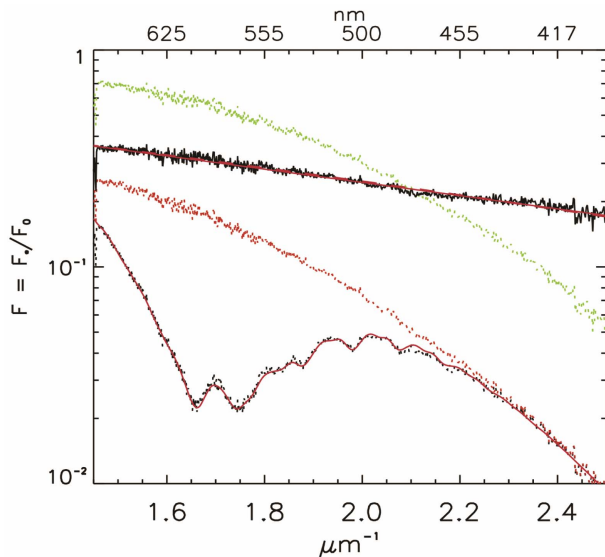
In 1871 Rayleigh [1] first attributed the blue of the sky to scattering of sunlight by small (compared to the visible wavelengths) particles and showed that the scattering law must vary as  $1/\lambda^4$ . Rayleigh also established that the irradiance of scattered (by small particles) and direct lights from the sun should be in the ratio

$$\frac{I}{I_0} \propto R = \frac{1}{\lambda^4} e^{-\frac{a}{\lambda^4}} \quad (2)$$

In 1899 Rayleigh finally concluded the small particles were air molecules [2]. In contradistinction with aerosol scattering, Rayleigh scattering is quasi-isotropic. Coefficient  $a$  in Equations (1) and (2) can be converted to nitrogen column density by  $N_{N_2} / a \approx 2.45 \cdot 10^{35} \text{ cm}^{-2}$  [23].

Hartley, and Chappuis to a lesser extent, suggested an alternative to Rayleigh's theory, absorption by ozone. Chappuis in his 1882 memoir [6] identifies a quasi-continuum of ozone absorption between 500 nm and 660 nm, now known as the Chappuis bands (**Figure 1**). He notes that "la couleur bleue caractérise l'ozone aussi sûrement que son odeur... La couleur que possède alors le gaz rappelle la couleur bleue du ciel"<sup>4</sup>. Low resolution spectroscopic observations by Wulf, Moore, and Melvin in 1934 [17] confirmed the presence of the Chappuis bands in blue twilight skies but not, or too weakly to be detected, in blue daylight.

The occultation spectrum of **Figure 1** was observed under a much larger air-mass than sun-rays experience when the sun is above the horizon. In day-time and in the visible wavelength range the extinction of sunlight in the atmosphere is much weaker: the Chappuis bands, with column densities between  $5 \cdot 10^{18} \text{ cm}^{-2}$  and a few  $10^{19} \text{ cm}^{-2}$ ,



**Figure 1.** A GOMOS (Global Ozone Monitoring by Occultation of Stars) occultation spectrum of star Sirius at a tangent altitude of 14780 m (courtesy of J.-P. Bertaux and L. Blanot, Service d'Aéronomie, Verrières-le-Buisson, France). Bottom dotted spectrum is the ratio  $F_{\star}/F_0$  ( $\propto T$  in Equation (1)) of the spectrum of Sirius observed through and outside of the atmosphere. The red plain curve is the fit (ozone + Rayleigh + aerosol extinction) of the occultation. Aerosol extinction ( $e^{-\frac{b}{\lambda^p}}$ ,  $p \approx 1.3$ ,  $b \approx 0.45$ ) changes (slightly) the slope of the spectrum (upper black trace). The green dotted trace shows the effect of Rayleigh extinction alone ( $N_{N_2} \approx 2.10^{34} \text{ cm}^{-2}$ ,  $a = 0.075$ ). The dotted red trace is  $F_{\star}/F_0$  corrected for ozone absorption ( $N_{O_3} = 3.9 \times 10^{20} \text{ cm}^{-2} \equiv 1.4 \times 10^4 \text{ D.U.}$ ).

<sup>4</sup>"Blue color characterizes ozone as certainly as its odor... The color of the gas then reminds one of the blue color of the sky."

<sup>5</sup>Typical zenithal ozone column densities are of order  $10^{19} \text{ cm}^{-2}$ , 370 D.U.

are imperceptible<sup>5</sup>; the modification of the spectrum's slope by Rayleigh extinction is also small,  $N_{N_2}$  must not exceed a few  $10^{32} \text{ cm}^{-2}$ .

It is only when the sun is close to the horizon that ozone absorption is large enough to produce the large depression between 1.4 and 2.2  $\mu\text{m}^{-1}$  on **Figure 1**. The column density of ozone along the optical path of sun-rays can rise to a few  $10^{20} \text{ cm}^{-2}$  [20]. The other important

effect is Rayleigh extinction ( $e^{-\frac{a}{\lambda^4}}$  in Equation (1)) which bends the blue side of the spectrum (green trace of **Figure 1**). Aerosol extinction acts as a lever and changes the average slope of the spectrum (black exponential trace on **Figure 1**). This change is however marginal in that it does not modify, as shown by the comparison of the green and red-dotted traces of **Figure 1**, the spectrum of the sun as much as Rayleigh extinction or ozone absorption do. In practice and from consideration of a spectrum alone aerosol extinction is difficult to disentangle from Rayleigh extinction unless it is large enough to modify the slope of the spectrum at low wave-numbers.

#### 4. Blue Sky

**Figure 2**, left frame, plots spectra of different types of blue skies: three blue day-time spectra (traces (1)) in different directions observed in the mid-afternoon (3~P.M. local time, observations of August 28, 2010), the sun is then well above the horizon; and two twilight blue sky spectra, one on the horizon  $180^\circ$  from the sun and one at the zenith observed 5 minutes later (observations of October 11, 2010). As found by Wulf, Moore, and Melvin [17] the major difference between daylight and twilight blue skies is the presence of ozone in the latter spectra.

Analytically the blue sky spectra of **Figure 2** deduce one from the other by a  $T$ -transformation of Equation (1) (**Figure 2**, right plot), and differ by the amount of Rayleigh extinction, ozone absorption, and aerosol extinction, sun-rays have experienced on their way through the atmosphere. Zenith spectrum (1) (upper trace) is well fitted (yellow spectrum of **Figure 2**) by the spectrum of the sun<sup>6</sup> times a Rayleigh function  $R$  (Equation (2)) with exponent  $a = 0.005$  ( $N_{N_2} = 1.2 \cdot 10^{33} \text{ cm}^{-2}$ ). The generic analytical expression of a blue sky spectrum is therefore

$$\frac{1}{\lambda^4} TS_0.$$

Ozone and nitrogen column densities for each spectrum are given in **Table 1**. Aerosol extinction could not be separated from Rayleigh extinction and aerosol column densities are omitted. Ozone is undetectable in the zenith daylight spectrum and must be less than  $10^{19} \text{ cm}^{-2}$ . The two following daylight spectra observed in the counter direction of the sun at about  $45^\circ$  and  $60^\circ$  from the

<sup>6</sup>The spectrum of the sun was provided by G. Thuillier [26].

zenith have comparable nitrogen column densities and a light ozone absorption ( $\approx 2.10^{19} \text{ cm}^{-2}$ , 750 D.U.). For the horizon and zenith twilight spectra ozone column densities are about a factor of ten higher.

## 5. The Red Color of the Sky

Spectra (4)-(6), of **Figure 3** are sunrise spectra observed

in the reddest region of the horizon. There is an average of 20 minutes interval between the three observations, starting with spectrum (6), when sunlight first appears. The sun is below the horizon. The steepness of the spectra, given by the relative weight of their red and blue parts, decreases as sun rises and is due to Rayleigh extinction. The simultaneous diminution of ozone absorption in the

**Table 1. Spectra characteristics (ozone absorption, rayleigh extinction).**

| $n^{(a)}$          | Sky type                   | $N_{\text{O}_3}^{(b)}$    |                     | $N_{\text{N}_2}^{(c)}$    | $R^{(d)}$       |                            |                        |
|--------------------|----------------------------|---------------------------|---------------------|---------------------------|-----------------|----------------------------|------------------------|
|                    |                            | $10^{19} \text{ cm}^{-2}$ | $10^3 \text{ D.U.}$ | $10^{32} \text{ cm}^{-2}$ | $a$             | $(1/\lambda)_{\text{max}}$ | $\lambda_{\text{max}}$ |
|                    |                            |                           |                     |                           | $\mu\text{m}^4$ | $\mu\text{m}^{-1}$         | nm                     |
| (1)                | Blue day-time              | 1                         | 4                   | 12                        | 0.005           | 3.76                       | 270                    |
| (2)                | Blue horizon (twilight)    | 16                        | 60                  | 130                       | 0.053           | 2.08                       | 480                    |
| (3)                | Blue zenith (twilight)     | 14                        | 52                  | 44                        | 0.018           | 2.73                       | 370                    |
| (4) <sup>(e)</sup> | Pink horizon (twilight)    | 18                        | 66                  | 320                       | 0.130           | 1.66                       | 600                    |
| (5) <sup>(e)</sup> | Reddish horizon (twilight) | 30                        | 110                 | 460                       | 0.190           | 1.51                       | 660                    |
| (6) <sup>(e)</sup> | Red horizon (twilight)     | 34                        | 126                 | 500                       | 0.200           | 1.46                       | 670                    |

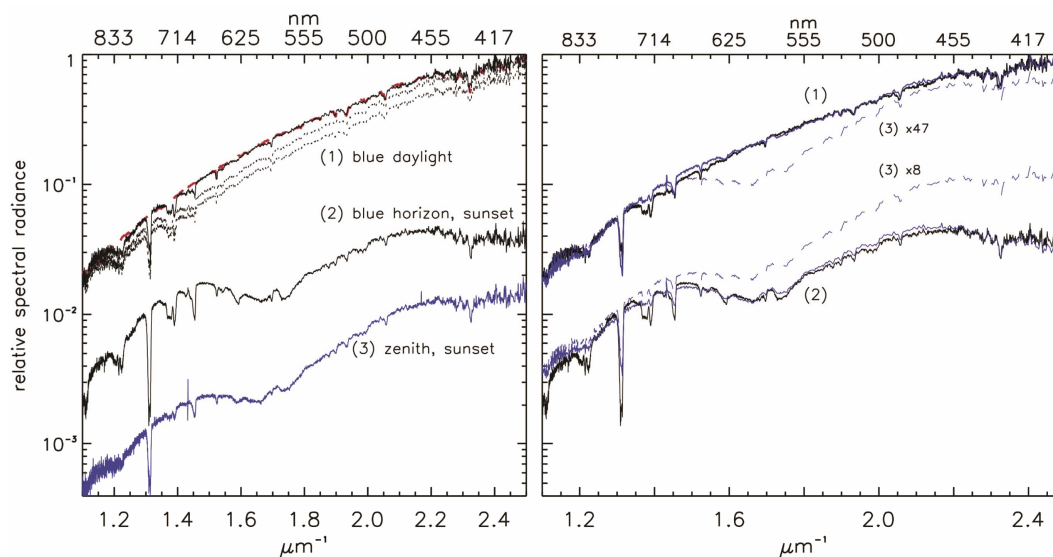
(a) Spectrum number in the figures.

(b) Approximate ozone column density along the path of sunrays from outside the atmosphere to the observer.

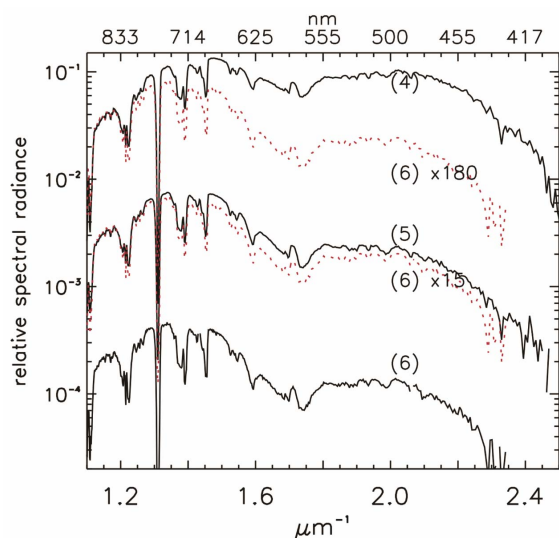
(c) Approximate nitrogen column density along the path of sunrays from outside the atmosphere to the observer.

(d) Parameters of the rayleigh function  $R$  (Equation (2)): value of  $a$ , of  $1/\lambda$  and  $\lambda$  at maximum. FWHM are  $0.59(1/\lambda)_{\text{max}}$  and  $0.66\lambda_{\text{max}}$  for  $R(1/\lambda)$  and  $R(\lambda)$  respectively.

(e) Spectra (4)-(6) are part of a series of over fifty consecutive spectra observed between 6 and 7 A.M. on October 11, 2010. Exact times of observations are: 6:12, 6:29, 6:50 for spectra (6), (5), and (4). Exposure times are 60s, 10s, 0.2s.

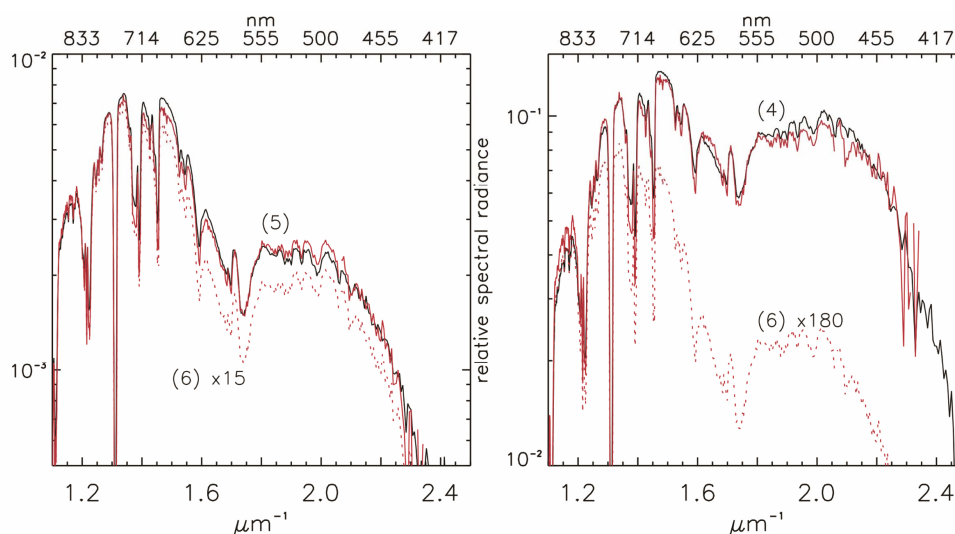


**Figure 2. Blue sky.** Left: three mid-afternoon blue sky spectra (top traces (1)), a twilight blue horizon in the direction opposite to the sun (2) and a zenith spectrum observed a few minutes later (3). The top trace is a zenith spectrum, the two dotted spectra were observed in the anti-solar direction at  $60^\circ$  and  $30^\circ$  elevation. The top red dashed line is a fit of the zenith blue sky spectrum by Rayleigh scattering of sunlight ( $\propto \frac{1}{\lambda^4} e^{-0.005/\lambda^4} S_0$ ). Right: blue spectra deduce one from the other (here spectra (1) and (2) are fitted by spectrum (3) corrected for ozone absorption and Rayleigh scattering, in blue) by a transformation  $T$ , Equation (1). The two dashed blue spectra are spectrum (3) which has been re-scaled. Ozone column densities for spectra (2) and (3) are of order (see Table 1)  $1.5 \times 10^{20} \text{ cm}^{-2}$  ( $5.5 \times 10^3 \text{ D.U.}$ ).



**Figure 3. Sunrise spectra.** Three sunrise spectra from very red (6) to pink (4). The two dotted red spectra are spectrum (6) which has been re-scaled to be compared to spectra (5) and (4).

Chappuis bands enlarges the red bump and moves its central wave-number towards the blue (right plot of **Figure 4**). In the same time the color of the horizon turned from red (spectrum 6) to pink (spectrum 4). The red color of the sky in these observations results from Rayleigh extinction which diminishes the blue part of the spectrum, and from ozone absorption which isolates the red bump on its blue side and shifts its central wavelength towards the red. As for blue skies all red horizons deduce one from the other by a  $T$ -transformation (**Figure 4**). The associated Rayleigh and ozone column densities along the path of light are given in **Table 1**.



**Figure 4. Red skies.** Spectra (4) and (5) deduce from spectrum (6) by a transformation  $T$ , Equation (1). Dotted red lines are spectrum (6) which has been re-scaled. Plain red traces are the fits of spectra (4) and (5) by spectrum (6) corrected for ozone absorption and Rayleigh extinction (see **Table 1**).

## 6. From Blue to Red Skies

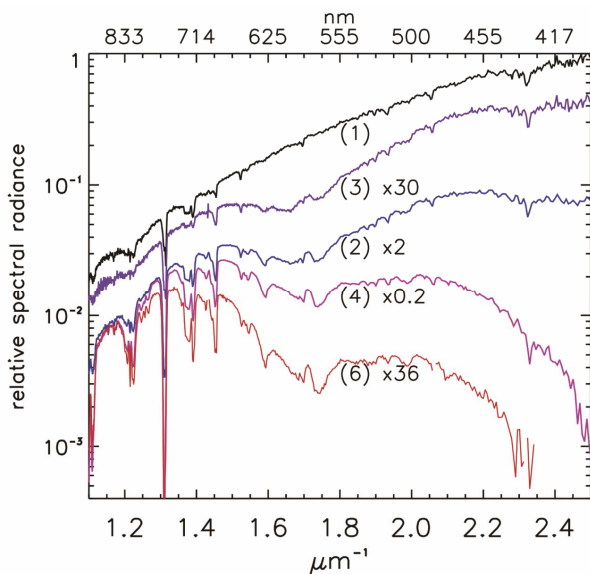
**Figure 5** is a summary of typical sky spectra that can be observed during a clear day. They are ordered by decreasing average slope and have been arbitrarily rescaled. The progressive change of color of the sky from blue daylight (spectrum (1)) to blue twilight (spectra (2) and (3)), pink (spectrum (4)), and red (spectrum (6)), is mainly due to Rayleigh extinction and ozone absorption.

The blue color of the sky in day-time (top spectrum (1) on **Figure 5**) is Rayleigh scattering by nitrogen of weakly attenuated sunlight. It is modified at twilight, when the sun is close to the horizon, because of ozone absorption which separates the red and blue sides of the spectrum (**Figure 5**, spectra (2) and (3)).

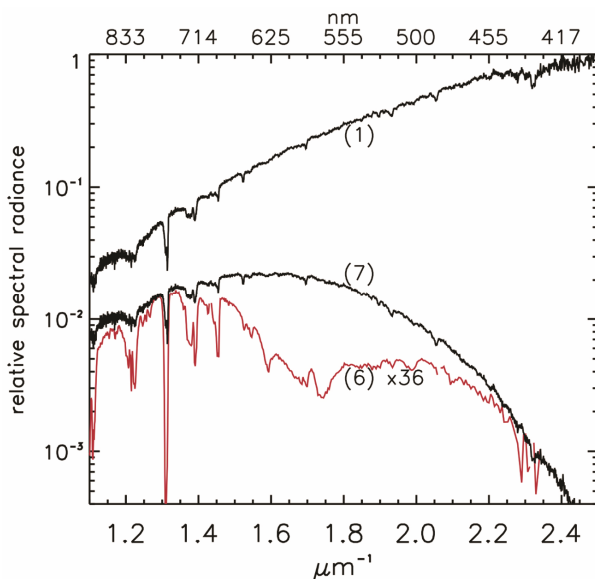
The much larger path of sunrays along the horizon at sunset or sunrise dramatically increases ozone absorption and Rayleigh extinction. The maximum of the Rayleigh function  $R$  (Equation (2)), which was in the UV or in the near-UV for blue skies (**Table 1**), is moved to the longer wavelengths (smaller wave-numbers). For spectrum (6) this maximum is close to  $1.65 \mu\text{m}^{-1}$  (600 nm, spectrum (7) in **Figure 6**). Without ozone the color of the sky would be that of spectrum (7), in the orange, as anticipated by Bohren and Fraser [16], and cannot alone explain the red color of the sky. Ozone provokes the large depression between spectra (6) and (7) which isolates the red bump centered at  $1.3 \mu\text{m}^{-1}$  (770 nm) and is responsible for the red color of the sky.

The ozone Chappuis bands absorb the green-orange part of the spectrum and have two roles in the color of the sky. For not too high optical depths they turn blue daylight into blue twilights, as it has generally been





**Figure 5. From blue to red skies. Spectra of blue daylight (1), blue twilight ((2) and (3)), pink (4), and red (6), colors of the sky. The spectra have been arbitrarily re-scaled and ordered by decreasing slope from top to bottom. They deduce one from the other by a transformation  $T$  (Equation (1)), and are  $\propto \frac{1}{\lambda^4} TS_0$ . The red bump moves to the left (thus to the red) as ozone column density and Rayleigh extinction increase.**



**Figure 6. Spectrum (7) is spectrum (1) (blue afternoon sky) with an additional Rayleigh extinction ( $N_{N_2} = 5 \times 10^{34} \text{ cm}^{-2}$ ). Absorption between spectra (7) and (6) in the Chappuis bands ( $N_{O_3} = 3.4 \times 10^{20} \text{ cm}^{-2} \equiv 12.6 \times 10^3 \text{ D.U.}$ ) gives the red color of the horizon.**

foreseen. But when Rayleigh extinction is large enough to inverse the slope of the spectrum in the blue, ozone

absorption also isolates the red bump of spectrum (6) and gives a major contribution to the red color of the sky.

## 7. Discussion

The purpose of these investigations was to provide a comprehensive understanding of the color of the sky. Spectral observations keep the signature of the media crossed by light. Spectroscopy of the sky at different times of the day shows that the relevant parameters that describe the color of a clear sky are the same that determine the extinction of sunlight: the exponents  $a$ ,  $b$ ,  $N_{O_3}$  which enter in the expression of the  $T$  function of Equation (1). The large variety of spectra that can be observed indicates that the color of the sky is far more nuanced than what the eye perceives. This panoply of colors is however limited by the presence of ozone which suppresses the orange-green part of the spectrum, when the sun is low on the horizon.

The Rayleigh function  $R$  of Equation (2) has a bump-like shape resembling that of spectrum (7) (Figure 6). As long as blue wavelengths lay on the low column density side (the  $1/\lambda^4$  low optical depth rise) of  $R$  the sky is blue, with nuances that depend on the mean slope of the spectrum (due to Rayleigh and aerosol extinction) and on the column density of ozone along the path of light. The sky turns to red colors when Rayleigh extinction becomes large enough to shift the maximum of  $R$  (Table 1) in the visible wavelength range and the blue wavelengths on

the exponentially (the  $e^{-\frac{a}{\lambda^4}}$  term in  $R$ ) decreasing side of the spectrum. The increase of Rayleigh extinction is accompanied by an increase of ozone absorption which isolates the red part of the spectrum and shifts its maximum to the red. Without ozone the sky on the horizon would take orange to green colors (depending on where the maximum of  $R$  falls). One must therefore conclude that ozone, because of the critical position of the Chappuis bands in between the blue and the red wavelengths, gives the final red color of the sky at sunset or sunrise.

The generic analytical expression of the spectrum of a clear daylight sky is  $\propto \frac{1}{\lambda^4} TS_0$ . Scattered sunlight received by an observer on earth can be considered as the result of two operations, extinction of sunlight across the atmosphere and single Rayleigh scattering. It may happen that two components are necessary to represent specific observations, for instance a blue component may need to be added to a red spectrum in transition directions between red and blue skies [19]. But a red horizon spectrum like spectrum (6) of Figure 5 is not the sum of two components as its appearance might suggest.

Ozone absorption and Rayleigh extinction are the two major processes involved in the color of a sky illuminated by the sun. These observations did not highlight a significant role of aerosols in common daylight colors.

The aerosol extinction term  $e^{-\frac{b}{\lambda^p}}$  in Equation (1) is probably small and could not be disentangled, in these observations, from the change of slope due to Rayleigh extinction. It can only diminish the steepness of the  $1/\lambda^4$  rise of Rayleigh scattering, and the impression of red with it, which is not what is observed. But the analysis of sun occultation spectra [20] proves that before the sun completely sets its light is progressively replaced by sunlight forward-scattered by aerosols. If this light was in turn scattered by nitrogen in the direction of the observer, the wavelength dependence of the spectrum at small wave-numbers would be  $\propto 1/\lambda^5$  instead of  $1/\lambda^4$ . The red side of the spectrum would be steepened, and the impression of red increased. In a cloudy sunrise or sunset aerosol scattering that follows Rayleigh extinction would have the same effect. Therefore the possibility that aerosols contribute to the red color of the sky at the very beginning of a sunrise, in the latest part of a sunset, or under a cloudy sky, cannot be ruled out. Evidence of this phenomenon would require observations with a larger wavelength coverage of the near-infrared spectrum.

The analytical representation of the spectra and the way they transform one into the other excludes any significant participation of multiple scattering in the color of the sky. This is not surprising given that Rayleigh scattering is, in contradistinction with forward scattering by aerosols, nearly isotropic. The loss of photons over all directions after each scattering renders multiple Rayleigh scattering a particularly inefficient process.

Molecular absorption by molecules other than ozone consists in local interruptions of the spectra, is marginal except for the strong water bands, and was neglected. Water vapors are important at sunrise (spectra (4)-(6)) but do not appear to influence significantly the color of the sky.

These observations could be improved in several ways, especially by extending the wavelength coverage in the infrared. By combining photography with spectral observations of the sky and a more systematic planning of observations, the link between spectra and the palette of colors of the sky would be better captured. Finally a long lasting and controversial issue, the green-flash, which is also observed when the sun is low on the horizon, could hopefully be solved with spectroscopy. All these issues suggest interesting research projects that can be studied from the college level on with a relatively inexpensive equipment.

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